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Study on irreversible degradation mechanism of organic-inorganic hybrid perovskite materials by trapped charges

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서울대학교 대학원
기계항공공학부
곽 귀 성
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Study on irreversible degradation mechanism of organic-inorganic hybrid perovskite materials by trapped charges

Kwisung Kwak
School of Mechanical and Aerospace Engineering
The Graduate School
Seoul National University

Abstract

Organic-inorganic hybrid perovskite crystal has excellent photo-electronic properties, so it is used as various photo-electronic devices. In particular, organic-inorganic hybrid perovskite solar cells are expected to replace silicon solar cells because of their high efficiency and low production costs. The biggest problem of commercialization of organic-inorganic perovskite solar cell is low stability. There are various reasons for the lack of stability of the perovskite solar cell, but the most important reason is that the perovskite material is easily decomposed when light is irradiated in air. In addition, perovskite material is degraded by various factors. Therefore, it is
essential to clarify the origin and decomposition process of irreversible
degradation of perovskite materials to ensure long-term stability of
perovskite solar cells.

First, the degradation of CH$_3$NH$_3$PbI$_3$ (MAPbI$_3$) material, which is most
commonly used in perovskite solar cells, was studied. In the case of
perovskite solar cells, efficiency, hysteresis, and degradation patterns were
different depending on the electron transport layer. It is assumed that the
origin of the degradation of the perovskite materials is charge because it is
a characteristic occurred by the behavior of charges. To demonstrate the
effect of charges, a corona discharge method was used to inject charges
without light irradiation. The transfer of the charges injected by the corona
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spectroscopy. The perovskite material was decomposed under the
conditions of the injected charge and water or oxygen. These results
demonstrate that the trapped charges in the perovskite are the origin of the
decomposition of the perovskite material. In order to analyze the theoretical
decomposition mechanism of perovskite materials, the ab initio molecular
dynamics (AIMD) simulation based on density functional theory (DFT)
was carried out. As a result of the simulation, when the trapped charges
were present, it was observed that the perovskite crystal collapsed by
reacting with moisture or oxygen, but there was nothing happened without
the trapped charges.

Second, the investigation of different composition of perovskite material
and its degradation mechanism were studied. The crystal structure can be
changed from tetragonal structure to a cubic structure by changing the organic cation or halide anion in MAPbI$_3$ through the calculation of the goldschmidt tolerance factor. In the case of the cubic structure, the structural stability is superior to the tetragonal structure and high efficiency can be obtained. Therefore, the perovskite composition which has the cubic structure was investigated using CH($(\text{NH}_2)_2$)$^+$ (FA$^+$) organic cation instead of CH$_3$NH$_3^+$ (MA$^+$) organic cation and Br$^-$ halide anion instead of I$^-$ halide anion. As a result, we found the composition of MA$_{0.6}$FA$_{0.4}$PbI$_{2.9}$Br$_{0.1}$, which has an efficiency of 20.2% and superior stability to MAPbI$_3$. This composition also was decomposed in the presence of trapped charge, but it showed high stability in oxygen atmosphere compared with MAPbI$_3$.

Recently, FAPbI$_3$ material, which has cubic structure, has been widely used because of its high light absorption properties and excellent light stability. However, since the stability of FAPbI$_3$ material is also not sufficient for commercialization, the degradation mechanism of FAPbI$_3$ material was studied. In the case of FAPbI$_3$, it was confirmed that it rapidly changed from a cubic structure to a hexagonal structure under high relative humidity condition. However, when oxygen was injected under the light irradiation condition, it was confirmed that the crystal was decomposed into PbI$_2$. This process was also confirmed by the corona discharge method and the AIMD simulation.
Keywords: perovskite solar cell, organic-inorganic hybrid perovskite, trapped charge, electron transport layer, degradation mechanism, corona discharge, AIMD simulation, cubic structure, mixed compositional perovskite,

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Nomenclature

\( V_{oc} \)  out-circuit voltage \([\text{V}]\)

\( J_{sc} \)  short-circuit current density \([\text{mA cm}^{-2}]\)

\( U_{\text{offset}} \)  Offset voltage for suppression charge extraction \([\text{V}]\)

\( U_{\text{max}} \)  Maximum voltage \([\text{V}]\)

\( j_0 \)  Capacitive displacement current \([\text{mA cm}^{-2}]\)

\( A \)  Voltage increase rate \([\text{V/s}]\)

\( t_{\text{max}} \)  Time at maximum current \([\text{s}]\)

\( \mu \)  Carrier mobility \([\text{cm}^2 \text{V}^{-1} \text{s}^{-1}]\)

\( r_A \)  Radius of A site cation \([\text{nm}]\)

\( r_x \)  Radius of X site anion \([\text{nm}]\)

\( r_{\text{Pb}} \)  Radius of Pb\(^{2+}\) \([\text{nm}]\)

\( t \)  Goldschmidt tolerance factor
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Chapter 1. Introduction

1.1. Research Background

Research on renewable energy for replacing petroleum energy is very active worldwide. Although there are many types of renewable energy, the demand for solar cells using infinite solar energy is continuously increasing.[1] Currently, the solar cell industry is dominated by Si-based solar cells. However, these are many disadvantages in Si-based solar cell, for example expensive production cost, so many researchers are studying various next-generation solar cells which can replace Si-based solar cells.

Among the next generation solar cells, perovskite solar cells have recently been shown to be the most promising solar cells to replace Si-based solar cells, with a remarkably high efficiency of more than 23%. These perovskite solar cells have shown very rapid efficiency growth since their first reported in a group of prof. Miyasaka in 2009.[2] CH$_3$NH$_3$PbI$_3$ (MAPbI$_3$) material is the most widely used and studied material in the study of perovskite solar cells. MAPbI$_3$-based perovskite solar cells has been studied very actively because it is easy to manufacture and shows about 20% efficiency.[3] In addition, MAPbI$_3$-based perovskite solar cells can be processed at
under 150°C, and they are still widely used in various photo-electronic devices such as flexible devices.[4] In recent years, in order to achieve higher efficiency, mixed composition perovskite solar cells using a mixture of various organic cations instead of MAPbI$_3$ have been studied.[5-7] Among these various composition perovskite solar cells, CH$_2$((NH)$_2$)$_2$PbI$_3$(FAPbI$_3$)-based perovskite solar cells are also vigorously studied because they show high power conversion efficiency and excellent light and thermal stability. [8,9]
Figure 1.1 Schematic image of perovskite crystal structure adopted from [10].
1.1.1. The stability of MAPbI₃ perovskite materials.

Despite of rapid efficiency growth, perovskite solar cells still fail to commercialize because the stability of perovskite is very poor. Although perovskite solar cells have been reported to have stability over 1000 hours in recent years,[11] they have not yet achieved sufficient stability for commercialization. There are many reasons why the efficiency of perovskite solar cells is diminishing. Many studies have been reported that the stability of the charge transport layer constituting the perovskite solar cell and the migration of metal electrode affect the stability of perovskite solar cells.[12-14] But the biggest problem is that the stability of perovskite material itself is poor.[15-17] In particular, the stability of MAPbI₃ material, which is the most frequently used, is very terrible, and therefore, the efficiency of a MAPbI₃-based perovskite solar cells is widely known to decrease rapidly in a short time.[17,18] The MAPbI₃ material is easily transformed into PbI₂ due to the collapse of the crystal structure.[17]

There were studies to investigate the causes of degradation of MAPbI₃ perovskite materials as seen in Figure 1.1. First, there is a reversible hydration state due to moisture.[15,19] In the case of the MAPbI₃ material, it becomes a hydration state(CH₃NH₃PbI₃·H₂O or (CH₃NH₃)₄PbI₆·2H₂O) in combination with moisture under a high
relative humidity environment. The hydration state of MAPbI$_3$ has a transparent color, which means that it cannot absorb the light. Therefore, the hydration state of MAPbI$_3$ materials does not function as a light absorber, and thus the solar cell is not driven. This transparent hydration state return to its original black perovskite state in the absence of moisture, but it does not act as a solar cell during the hydration state.

Next, photo-degradation by photocatalytic effect under light illumination was reported in a group of prof. Ito.[20] When the electron transport layer is made of TiO$_2$, TiO$_2$ absorbs light under light illumination condition, and the photocatalytic effect of TiO$_2$ destroys MAPbI$_3$ into PbI$_2$. Recently, perovskite solar cells using different types of electron transport layers have been reported to replace TiO$_2$.[21,22]

The degradation by strong electric field was by a group of Petrozza and Snaith.[23] After deposition of two Au electrodes on the perovskite surface, the application of a strong electric field between the two electrodes causes degradation from near the electrodes. At this time, the electric field applied between the two electrodes is about 600 V / cm, and they said that the strong electric field causes the migration of organic cations and finally the perovskite crystals collapse.
The group of prof. Haque reported that degradation of MAPbI$_3$ material occurs when light is irradiated in an oxygen atmosphere.[16,24,25] When light is irradiated on MAPbI$_3$, excited electrons are generated. As the electron is transferred to oxygen, the oxygen becomes a superoxide state, and this superoxide decomposes MAPbI$_3$ material. In particular, during the MAPbI$_3$ degradation process, they speculated the deprotonation of organic cation. The hydrogen atoms in MAPbI$_3$ are separated and combine with oxygen to form water, and MAPbI$_3$ decomposed into PbI$_2$.

In addition, a variety of causes have been identified as the origin of the degradation of MAPbI$_3$. However, it is unclear what process MAPbI$_3$ material will undergo during the degradation.
Figure 1.2 (a) Schematic of reversible hydration formation of MAPbI$_3$ material adopted from [15]. (b) Degradation of MAPbI$_3$ material by photo catalytic effect of TiO$_2$ adopted from [20] (c) Ion migration driven degradation of MAPbI$_3$ material by electric field adopted from [23] (d) Decomposition of MAPbI$_3$ material by super oxide generated under oxygen ambient adopted from [24].
1.1.2. The stability of FAPbI$_3$ perovskite materials

Recently, the perovskite solar cells using black alpha phase FAPbI$_3$ (α-FAPbI$_3$) material as a light absorber have been reported in recent years.[8,9] α-FAPbI$_3$ material has higher light absorption efficiency and light and thermal stability than MAPbI$_3$. However, pureα-FAPbI$_3$ material is unstable in the presence of water at room temperature. Because the Goldschmidt tolerance factor of pure FAPbI$_3$ is greater than 1 at room temperature, this means that the hexagonal delta phase FAPbI$_3$(δ -FAPbI$_3$) is more favorable than cubic alpha phase FAPbI$_3$. (The structure of FAPbI$_3$ shown in Figure 1.3) Especially, α-FAPbI$_3$ is very vulnerable to moisture, so it changes very easily into δ -FAPbI$_3$ in the presence of water. The δ -FAPbI$_3$ material is yellowish, so it cannot act as a light absorber. Recently, it has been reported that α -FAPbI$_3$ material is maintained more stable in the presence of water at room temperature when a small amount of Cs cation enter the FAPbI$_3$ crystal structure.[8]
Figure 1.3 Schematic illustrations of two phases of FAPbI$_3$ material adopted from [26]. (a) $\alpha$-FAPbI$_3$ have the cubic perovskite structure. (b) $\delta$ -FAPbI$_3$ have 1-Dimensional hexagonal structure.
1.1.3. Enhanced stability achieved by changing the charge extraction layer

It is well known that the efficiency and stability of perovskite solar cells depend on the charge extraction layer. Therefore, it is very important to use an appropriate charge extraction layer for the perovskite solar cells.

The prof. Han’s group reported the perovskite solar cell using Nb doped TiO$_x$ electron transporting layer. (shown in Figure 1.4 a) In this device, the rapid electron extraction were achieved by fabrication of 10-20nm of Nb doped TiO$_x$ on the PCBM. As the result, a solar cell having no hysteresis and having higher efficiency and stability was fabricated.

In prof. Seok’s group, La-doped BaSnO$_3$ was used as an electron transporting layer instead of TiO$_2$. (shown in Figure 1.4 b) In case of BaSnO$_3$, it has excellent electrical mobility, so charge extraction is easy and therefore it shows high stability.

In prof. Yang’s Group, all charge extraction layers were made of inorganic material to improve stability. In this case, the charge extraction efficiency was much lower than that of the organic material, but the stability was much better. This is because the use of an inorganic transport layer prevents penetration of the air into the perovskite layer, thereby preventing the perovskite from being
decomposed.

Two factors play an important role in the stability of the perovskite solar cell. The first is to efficiently extract the charge from the device. The higher the extraction efficiency of the charge, the more stable the device is. Secondly, it prevents the infiltration of the air. It is very important to shut off this outside air because the air, especially moisture or oxygen, has a direct effect on the perovskite material. From these results, it can be predicted that an important factor in the decomposition phenomenon of perovskite materials is the reaction of charge and moisture or oxygen.
Figure 1.4 (a) Schematic illustration of the device configuration using the doped charge carrier extraction layers adopted from [27]. (b) Long-term stability of perovskite solar cells using doped electron extraction layer adopted from [28].
Figure 1.5 Comparison of the stability of devices based on inorganic charge extraction layer and organic charge extraction layer adopted from [29].
1.2. **Research objectives**

The dissertation aims to clarify the causes of the degradation of perovskite materials more clearly.

As seen in Table 1.1, when light was illuminated in the presence of moisture or oxygen, and electric field was applied to the perovskite films between gold two electrodes in humid air, the irreversible degradation of perovskite materials occurred. Light illumination and electric field have one common point. Those are associated with charges. In other words, charges can be generated by light illumination and an electrical current through perovskites. From these results, it can be inferred that charges could be the key cause that initiates an irreversible degradation of perovskite materials. To demonstrate this speculation, a novel experimental setup was introduced and a density functional theory (DFT)-based ab initio molecular dynamics (AIMD) simulation was performed to predict a more theoretically accurate mechanism.

First, the degradation of MAPbI$_3$ material was studied. It was found that trapped charges are the key cause of the degradation of the perovskite solar cells. When the electron transport layer was changed from TiO$_2$ to C$_{60}$, the degradation rate and degradation pattern were completely different. It was also confirmed that the hysteresis of the devices depends on the type of electron transport...
layer, which means that there are accumulated charges in the perovskite layer. To demonstrate this trapped charge driven mechanism, a novel experimental device for intentionally trapping charges on the surface of perovskite was introduced. A corona discharge method was used to inject charges on the surface of perovskite layer without light irradiation. It was confirmed that the charges injected by corona discharge method or charges generated by the light are collected at the grain boundaries and some spot of grain surface by Kelvin Probe Force Microscopy (KPFM) measurement. It was also confirmed that the decomposition of perovskite material occurs when there are trapped charges in the presence of water or oxygen. The AIMD simulation based on the DFT calculation was used to clarify the degradation mechanism by the trapped charge.

Second, the search and the degradation mechanism of perovskite materials other than MAPbI$_3$ were studied. MAPbI$_3$ has a tetragonal structure so it is structurally unstable. To overcome this structure instability, it is necessary to search for perovskite materials having a cubic structure. As the Goldschmidt tolerance factor approaches 1, the perovskite crystal structure becomes closer to the cubic structure. In the case of MAPbI$_3$, since the size of CH$_3$NH$_3$\textsuperscript{+}(MA\textsuperscript{+}) cation is small, it has tetragonal structure. If MA\textsuperscript{+} cation is replaced with CH((NH$_2$)$_2$\textsuperscript{+})(FA\textsuperscript{+}) cation, the perovskite crystal structure can be
close to cubic structure. As the result of stability and efficiency of various composition perovskite material, \( \text{MA}_{0.6}\text{FA}_{0.4}\text{PbI}_{2.9}\text{Br}_{0.1} \) showed best stability and showed 20.2% of power conversion efficiency. However, \( \text{MA}_{0.6}\text{FA}_{0.4}\text{PbI}_{2.9}\text{Br}_{0.1} \) material also was degraded in the presence of trapped charge and water molecule.

Finally, the degradation mechanism of\( \alpha \)-\( \text{FAPbI}_3 \) perovskite material was studied. In the case of \( \alpha \)-\( \text{FAPbI}_3 \), it was found that it easily changed into delta phase under high water conditions. However \( \alpha \)-\( \text{FAPbI}_3 \) material did not changed into delta phase but was decomposed into \( \text{PbI}_2 \). This degradation mechanism was also identified through AIMD simulation.
<table>
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<th>Light on</th>
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<th>Electric field</th>
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<tr>
<td>Nitrogen</td>
<td>No degradation</td>
<td>No degradation</td>
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<tr>
<td>Water</td>
<td>Irreversible degradation</td>
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<td>Oxygen</td>
<td>Irreversible degradation</td>
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Table 1.1 Summary of the causes of the degradation of perovskite material according to previous reports.
Chapter 2.

Trapped charge driven degradation of CH₃NH₃PbI₃(MAPbI₃) material.

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2.1. Introduction

Metal halide perovskites have been attracting worldwide interest and the power conversion efficiency (PCE) of perovskite based solar cells has already exceeded 22%. Their long-term stability issue is the most pressing problem for commercialization.[18] It is well known that perovskite materials are vulnerable to the exposure of humidity and light.[19,30,31] Though many efforts to encapsulate the devices for preventing direct contact to humidity have been attempted, it was not successful to obtain long-term stability comparable to commercial photovoltaic devices.[27] Various factors that can affect the stability have been investigated from the viewpoints of chemical structure[32], electrical stress[23], hydrated states[15,19,30] and heat[33]. However, the degradation mechanism is still unclear how a
fully fabricated device deteriorates rapidly even though its perovskite layer is tightly-covered by the hole transport material (HTM) and the back electrode. It is also elusive why light soaking causes irreversible degradation of perovskite materials in the presence of moisture while the moisture in the dark condition only induces reversible hydration of perovskite materials. [23,30]

The device structure and the choice of charge extraction materials also influence the stability. The use of inorganic charge extraction layers was reported to enhance stability. [11,27] Other studies have focused on eliminating the pathway of water vapor infiltration into the perovskite film by coating carbon-based materials, polymers, and hydrophobic materials on the top surface of the perovskite film. [31,34] These approaches provide the device lifespan longer than the case without these materials, but not long enough to ensure long-term stability. Such approaches also occasionally sacrifice the photovoltaic performance. Especially, the devices employing titanium dioxide (TiO₂) as electron transport layer (ETL) are rapidly degraded under light soaking [8] even though they are highly efficient in energy conversion and exhibit the world’s best efficiency.

In the present study, we demonstrate that the charges trapped at the interface between perovskite and charge extraction materials are responsible for the irreversible degradation due to moisture. The
elusive experimental puzzles both on the occurrence of degradation beginning from different side depending on different charge extraction material and the role of light soaking for irreversible degradation will be clearly explained from the present concept of the trapped charges later.

We also experimentally demonstrate that trapped charges also induce irreversible degradation of MAPbI$_3$ even in the absence of moisture as long as oxygen is supplied. These results suggest that the trapped charges lead to the degradation of perovskite material when there is oxygen or water or both. In our experiments, charges were generated either by light soaking or corona ion deposition[35], and trapped charges on the grain boundaries were confirmed by Kelvin Probe Force Microscopy (KPFM) measurements for either case of light soaking or ion deposition. To investigate the details of charge-driven degradation mechanism at the atomic level, we carried out AIMD simulations and tracked how oxygen and water molecules in the presence of trapped charges interact with the compositional units of MAPbI$_3$ crystal in the first 2 ps. From these calculations, we confirmed that deprotonation of organic cations could be induced by trapped charges in the presence of water. We also learned that oxygen molecules greatly weaken the MAPbI$_3$ crystal bonding when there are trapped charges even without water.
2.2. Results and discussion

First, we examined how solar cell degradation behavior becomes different depending on different charge extraction layers, for example, $C_{60}$ and compact TiO$_2$.

As shown in Figure 2.1a, b, non-encapsulated $C_{60}$ based cell shows much more stable performance under one sun illumination but still degrades while compact TiO$_2$ based non-encapsulated cell completely died only after 6 hours. To examine the detailed evolution of degradation, we investigated how the cross sectional morphology of the $C_{60}$ and TiO$_2$ based devices would evolve under illumination via the focused ion beam (FIB) assisted scanning electron microscope (SEM) images shown in Figure 2.1c, d. Consistent with the PCE measurement results, the SEM images clearly confirm that the $C_{60}$-based devices are much slowly degraded compared to the TiO$_2$-based cells. Strikingly, they showed different degradation patterns, namely, different side of degradation beginning where the degradation is initiated depending on different ETLs.

Since the reactants that can decompose perovskite materials could infiltrate from the thin metal electrode rather than from the thick ITO glass, it would be expected that the degradation should be initiated at the interface closer to the thin Au metal electrode. However, the perovskite films of TiO$_2$-based devices began to be decomposed at
the interface adjacent to the compact TiO$_2$ layer near FTO glass (as shown in Figure 2.1d). Although these observations were attributed to UV light induced photocatalytic effect of TiO$_2$ layer according to the previous report[20], we confirmed that the same degradation pattern happened in the TiO$_2$ and C$_{60}$ based devices aged even under UV-filtered light illumination (see Figure 2.2). For those of C$_{60}$-based devices, the decomposition began from the interface adjacent to HTL near Au metal electrode opposite to the case of TiO$_2$ based devices. Since the two types of devices have identical structure except for the ETL, C$_{60}$ or TiO$_2$/Perovskite/Spiro-MeOTAD/Au, these different degradation characteristics indicate that charge extraction may play an important role where moisture driven decomposition of perovskite material begins.
Figure 2.1 (a,b) Time evolution of the normalized PCE measured under one sun illumination in ambient conditions (relative humidity = 30\%) of the (a) C60 and (b) TiO2 based devices. (c,d) Time evolution of the FIB-SEM cross-sectional images of the (c) C60 and (d) TiO2 based devices aged for 72 h under light illumination in ambient conditions. Scale bars = 200 nm.
Figure 2.2 Time evolution of the FIB-SEM cross-sectional images of the C\textsubscript{60} (left) and TiO\textsubscript{2} (right) based devices, which were aged for 48 h under UV filtered light illumination in ambient conditions. Scale bars = 200 nm.
Second, the TiO$_2$ based devices suffer from severe hysteresis, whereas the C$_{60}$ based devices do not. Considering that the origin of hysteresis is known as capacitive current[36], trapped charge[37] and unbalanced charge injection[38], many electrons may be accumulated near the ETL in the TiO$_2$ based devices, while the C$_{60}$ based devices hardly do. From the observation on the degradation of the TiO$_2$ based devices that begins from the interface contacting TiO$_2$ layer where many charges could be trapped, it is reasonable to suspect that trapped charges at the interface between perovskite and charge extraction layer would be responsible for initiating the moisture related decomposition. Fast extraction of electrons through C$_{60}$ would hardly accumulate negative charges at the interface between perovskite and C$_{60}$, but hole extraction through Spiro-MeOTAD could be slower than the rate of electron extraction in the C$_{60}$ based cell.[39,40] This could result in positive charge trapping at the interface between perovskite and hole extraction layer, which could be the cause why the degradation begins from the interface between perovskite and Spiro-MeOTAD for C$_{60}$ based cells (see Figure 2.1c). These results demonstrating the degradation beginning from opposite side for different charge extraction layers gave us a clue about the trapped charge driven degradation regardless of polarity.
Another intriguing experimental observation is the light soaking in the presence of moisture which consistently showed irreversible degradation of perovskite in previous works[18,30,31] while in the dark condition moisture introduction only formed reversible hydrates of perovskites, for example, CH$_3$NH$_3$PbI$_3$•H$_2$O or (CH$_3$NH$_3$)$_4$PbI$_6$•2H$_2$O.[23,30] The reason has not been elucidated yet although Christians et al.[30] suggested that organic cation could become less tightly bound to PbI$_6^{4-}$ octahedra after light soaking. In the present study, along with the scenario of the above mentioned trapped charges that could trigger irreversible degradation, the charge generation under light soaking and subsequent trapping on the surface of perovskite is suspected to initiate the moisture induced irreversible degradation under light illumination. To confirm the irreversible degradation under light soaking, we also investigated the experiments under light soaking or not in the presence of moisture. Figure 2.3a showed the degradation behavior of MAPbI$_3$ for two days in the dark condition with relative humidity (RH) 90 %. Absorption spectra measurements show that the original MAPbI$_3$ (black curve) became hydrated (red curve) after two days and then dehydrated reversibly via N$_2$ drying, which is consistent with previous studies.[23,30] However, MAPbI$_3$ showed irreversible degradation rapidly under light soaking even at low RH 20 % (see
Figure 2.3b), which is also consistent with previous report.[30,31,39]

As mentioned earlier, perovskite absorbing light can generate and store charges due to its capacitive property[41] that may be trapped on the grain boundaries (we will show trapped charges along grain boundaries after light soaking later). It could be hypothesized that these trapped charges generated under light soaking would be responsible for irreversible degradation, which is in line with the aforementioned hypothesis of trapped charge driven degradation explaining the initiation of degradation on different side depending on different charge extraction layers.
Figure 2.3 (a) Absorption spectra of MAPbI$_3$ perovskite films under dark conditions at 90% relative humidity. MAPbI$_3$ perovskites were transformed into transparent hydrated states after 2 days. (b) Time evolution of absorption spectra of MAPbI$_3$ during light soaking at 20% relative humidity. Pictures of the perovskite films before and after aging are shown on the right side of each figure.
To prove this compelling hypothesis of trapped charge driven degradation, we have configured a novel experimental setup employing an ion generator by corona discharge and a stainless chamber that blocks all incident light from outside as shown in Figure 2.4a. The air inside the chamber is isolated from the outside and controlled by two gas inlets that are connected to the independent gas sources (Gas 1 and Gas 2) (see more details in Experimental methods: Experimental setup for ion generation and deposition). Gas 1 is ionized by applying a high voltage to the pin of the corona chamber, delivered to the deposition chamber by gas flow, and deposited on the perovskite film placed at the bottom of the deposition chamber. Gas 2 passes through a water bubbler to regulate the humidity level in the deposition chamber. We measured the time evolution of the absorption spectra as the perovskite films were gradually degraded in the deposition chamber. Gas 1 was chosen as nitrogen or hydrogen for generating positive or negative ions for being used as different polarity charges trapped on the surface of perovskite, respectively while Gas 2 was nitrogen (see more details in Experimental methods: Corona ion generation). First, we needed to check that N$_2$ positive ions and H$_2$ negative ions themselves do not affect the degradation without moisture (see Figure 2.5). Next, we examined the degradation behavior in the presence of moisture when
charged ions deposited on the surface of perovskite. When the positively charged N\textsubscript{2} ions were deposited and the relative humidity in the chamber was held at 40\%, the perovskite film rapidly decays as shown in Figure 2.4b1. The deposition of negatively charged H\textsubscript{2} ions also showed the similar irreversible degradation behavior under the same moisture level in Figure 2.4b2. Although Figure 2.4b2 for negative charges appears to cause slower degradation than the case shown in Figure 2.4b1 for positive charges, this could not tell which polarity charges affect more adversely on the degradation since ion generation for different polarity in our experiment is different (see more details in Experimental methods: Corona ion generation). This suggests that the irreversible degradation of perovskite materials only take place when both moisture and charges exist simultaneously. Based on these observations, the degradation mechanism could be thought of two-step process: the formation of hydrated perovskite by humidity and the irreversible decomposition by trapped charges. The first step of the formation of hydrated perovskite was already reported by several groups.[19,30] Here, we suggest that local electric field caused by trapped charge could distort electrostatically the structure of hydrated perovskite in which octahedral PbX\textsubscript{6}\textsuperscript{4-} interacts with both organic cation and H\textsubscript{2}O and trigger the initiation of irreversible decomposition of perovskite (will discuss more
Leijtens et al. [23] found the irreversible degradation near the gold electrodes coated on perovskite film by applying a weak external field of 600 V cm\(^{-1}\) in the presence of moisture and attributed it to the ion movement through electric field. Since an electric field was applied between two electrodes touching perovskite film, electric current could flow and there was a possibility of charge trapping underneath the electrode, which might have played a role for degradation. To isolate the effect of pure external electric field, we examined the degradation of perovskite materials by applying non-contact electric field which was given by two floating electrodes; one electrode is in air above perovskite film coated on ITO glass and the other electrode exists beneath the glass. We found no degradation up to 12 kV cm\(^{-1}\) under 90% RH (See Figure 2.6). Note that this field will be dropped across the air gap and therefore, the real field inside perovskite film should be different from the given field and the perovskite film might be uniformly polarized by one-directional strong E-field because perovskite materials have a high dielectric constant. [42,43] Therefore, further study should be needed to completely understand the effect of pure external electric field. It is noted that the differences between the fields due to the trapped charges and the external field lie in the point-like character of the trapped charges, which produce locally
huge and irregular fields. The huge and irregular fields formed by charges trapped along grain boundaries could help the process of deprotonation.[20,44]
Figure 2.4 (a) Experimental setup of corona discharge for ion generation, bubbling system for humidification, and SUS chamber for ion deposition and blocking light. (b) Absorption spectra of the perovskite film measured at an interval of 18 hours during deposition of (b1) positive nitrogen ions and (b2) negative hydrogen ions at 40% relative humidity.
Figure 2.5 Absorption spectra of the perovskite film under (a) continuous positive nitrogen and (b) negative hydrogen ion deposition in moisture-free dark condition.
Figure 2.6 The effect of non-contact high electric field on the degradation of the perovskite film at high relative humidity (90%) under dark condition.
We also employed three experimental apparatuses for our perovskite degradation studies. Using the first apparatus that has a sealed chamber (light blocked) with a gas inlet and outlet (Figure 2.7a), we can study the reaction between a gas and perovskite film without any external perturbation. The ambient gas in the chamber was adjusted by the different combination of inlet gases. The second apparatus (Figure 2.7b) is designed for intentional charge trapping in the perovskite film without light soaking by introducing N$_2^+$ corona ions (Same experimental setup as mentioned above). It is noted that gas molecules such as O$_2$ or H$_2$O are mixed with N$_2^+$ ions in the deposition chamber after N$_2^+$ corona ions generated in the corona chamber are injected into the deposition chamber. N$_2^+$ corona ions are selectively deposited on the perovskite surface by applied electric field in the deposition chamber. In the third apparatus (Figure 2.7c), we can study the degradation of perovskite film placed in a gas-controlled chamber under solar-simulated light to investigate the synergetic effect of light soaking and gas molecules.

First, to reconfirm the effect of water molecules on MAPbI$_3$ degradation in dark condition, pure nitrogen gas was humidified through a water bubbler and injected into the sealed chamber where a MAPbI$_3$ film was placed (at the relative humidity of 40% and flow rate of 2 lpm). There was no change in the absorption spectrum of the
MAPbI$_3$ film for 24 hours as shown in Figure 2.7d, which means that water molecules alone do not cause irreversible degradation of MAPbI$_3$ perovskites in dark condition, as was reported previously[17,45]. However, Figure 2.7e and f show that either ion charge deposition (without light) or light soaking causes irreversible degradation of MAPbI$_3$ film in 24 hours under the same humidified nitrogen gas. Both ion deposition and light soaking were shown to cause charge trapping along perovskite grain boundary, where degradation was initiated, and we proposed that trapped charge would be the main cause for irreversible degradation occurred under light soaking and moisture. However, it still remains unclear how water molecules interact with MAPbI$_3$ crystal in the presence of trapped charges.

Oxygen molecules have also been suggested as a cause of irreversible degradation in previous reports.[16,24,25,46] Haque and coworkers demonstrated that MAPbI$_3$ could be decomposed by light activated oxygen (superoxide O$_2^-$) even without moisture.[16,24,25] In their study, a degradation test was carried out under light soaking condition in the presence of oxygen. In the present study, we were able to examine the effect of oxygen molecules on the degradation of MAPbI$_3$ with either light soaking or ion charge deposition without light.
First, we found that injection of only dry air (N₂+O₂) without charge deposition or light irradiation did not cause any degradation (Figure 2.7g). On the other hand, the MAPbI₃ film underwent complete degradation in 24 hours under dry air with ion charge deposition (without light) (Figure 2.7h) or under light soaking (Figure 2.7i) as was the same in the presence of moisture (Figure 2.7e, f). In light of the fact that there was no degradation in nitrogen ambient gas (Figure 2.8), oxygen molecules are believed to be responsible for chemically interacting with the compositional units of MAPbI₃ crystal, as will be discussed in detail later.
Figure 2.7 a-c, Schematic illustration of experimental apparatuses for aging test, a, under dark condition, b, with N$_2^+$ deposition, and c, under one-sun irradiation. d-f, Absorption spectra of MAPbI$_3$ film measured before and after 1 day of aging with a continuous flow of humidified nitrogen gas, d, under dark condition, e, with nitrogen cation deposition, and f, under one-sun irradiation. g-i, Absorption spectra of MAPbI$_3$ film measured before and after 1 day of aging with dry air, g, under dark condition, h, with nitrogen cation deposition, and i, under one-sun irradiation.
Figure 2.8 Absorption spectra of perovskite films aged under one sun irradiation in pure nitrogen ambience.
Next we investigated how trapped charge could decompose perovskite material in time by examining morphology evolution via SEM analysis. As shown in Figure 2.9a (Tilted top-view and cross-sectional images) and Figure 2.10 (Top-view images), the degradation is initiated from the grain boundaries. As the reaction continues, the color of the film turns into yellow, indicating that the perovskite is irreversibly decomposed to PbI$_2$. It is interesting to dig into why degradation occurs from grain boundaries in line with our trapped charge mechanism. To check the distribution of trapped charges on the surface of perovskite after uniform ion deposition, we measured Kelvin Probe Force Microscopy (KPFM) of un-treated perovskite and ion-treated perovskite films (see more details in Experimental methods: Topography and Kelvin probe force microscopy). Figure 2.9b shows topology and surface potential distribution of the perovskite surface on which positive N$_2$ ions were uniformly showered. Striking coincidence between two images is the evidence that charges are preferentially trapped along grain boundaries. Overlapped image of topology and potential distribution shown in Figure 2.11 clearly demonstrates charges are trapped along grain boundaries even though ions are showered uniformly. For untreated sample, there is no correlation between topology and potential distribution (see Figure 2.11). With this charge trapping
along grain boundaries, experimentally observed degradation pattern following grain boundaries and the fact that the irreversible degradation occurs only when moisture or oxygen and charges exist together are the evidences that trapped charges would be responsible for the initiation of irreversible degradation of perovskite materials. It is now apparent that grain boundaries are the most vulnerable sites for the degradation because they provide charge accumulation sites as well as infiltration pathway of water vapor or oxygen.[19,25] Successful enhancement of stability utilizing high mobility inorganic charge extraction layers supports the present idea.[11,27] To further investigate the possibility that the above mentioned intriguing experimental observation of irreversible degradation under light soaking might be related to the mechanism of trapped charges of the present study, we measured KPFM images on the surface of perovskite after light illumination without ion deposition. As shown in Figure 2.11, charges are clearly trapped along grain boundaries for the sample soaked by light confirming that light soaking alone induces charge trapping along the grain boundaries of perovskite material like was done by introduction of ion charges in the dark condition. From our concept, these trapped charges can now trigger the irreversible degradation due to moisture or oxygen as the same happened when ion charges are deposited in the dark. Therefore, the
fundamental cause for irreversible degradation would be the same, that is, the trapped charges that could trigger the irreversible degradation under humid air. Such irreversible degradation under light soaking was reported in previous several reports[30,34], but, the reason has not been clearly elucidated so far although Christians et al.[30] suggested the lessened hydrogen bonding after photoexcitation as a possible cause. Here, we argue that trapped charges under moisture or oxygen would be responsible for the initiation of irreversible degradation under light or in the dark since light illumination always generates charges and traps the charges along grain boundaries as shown earlier. This can explain well why moisture itself without illumination or intentional ion deposition only hydrated perovskite reversibly. Light illumination under nitrogen gas without moisture and oxygen for one day was shown to hardly degrade the perovskite as shown in Figure2.8 which is strongly contrasted with the case of light illumination under humid air (see Figure 2.3b).

We also investigated charge trapping details through high resolution KPFM measurements (Figure 2.12). High-resolution KPFM images of the fresh samples under light soaking or with nitrogen ion deposition clearly show that the location of trapped charges is dominantly along grain boundaries and on some spots of
grain surfaces. We also found that oxygen-induced degradation took place along the grain boundaries and on grain surfaces where charges are mainly trapped, as shown by high-resolution KPFM and Field Emission-Scanning Electron Microscopy (FE-SEM) images of the fresh and degraded samples (Figure 2.13). Moreover, the evolution of surface potential distribution was observed during the whole degradation process, it was confirmed that charges continuously trapped along the grain boundaries (Figure 2.14). Since the morphology and area of the degraded sample changed with time, root mean squared roughness values of potential distributions increased during the degradation, which is indicative of occurring more widespread degradation with time (Figures 2.13 and 2.14). Thus, we may hypothesize that trapped charges along grain boundaries trigger irreversible degradation of MAPbI$_3$ perovskites in the presence of water or oxygen.

In addition, Photoluminescence(PL) spectroscopy measurements were performed to investigate charge trapping when positive N$_2$ ion deposited(Figure 2.15). In steady-state PL measurements, the detected PL intensity of MAPbI$_3$ film was decreased after positive N$_2$ ion deposition. In time-resolved PL(TRPL) measurements, the slope of PL decays was steeper after positive N$_2$ ion deposition. It is also confirmed that the first time constant($\tau_1$) achieved from TRPL by the
3-exponential fitting decreases from 45ns to 6.9ns after positive \( \text{N}_2 \) ion deposition. (Table 2.1) These results are similar to those obtained when measuring perovskite films with charge extraction layers.\cite{47}

In the case of PL measurements of the perovskite film with charge extraction layer, the charges generated by the light are extracted to the charge extraction layer, so the recombination of excitons is reduced. As the result, the PL intensity in steady-stat PL and first time constant in TRPL are decreased. Therefore, from the above results, it can be expected that the positive nitrogen ion deposited by the corona discharge method stays on the perovskite layer and acts as a trapped charges.
Figure 2.9 (a) Scheme description of perovskite degradation processes (left), and top-view (middle) and cross-sectional (right) SEM images of perovskite layers by ion deposition in humidified nitrogen. Scale bars = 200 nm. (b) Topography and surface potential profile of the perovskite film obtained from KPFM measurements after deposition of N₂ positive ions.
Figure 2.10 The top-view SEM images of the fresh and degraded perovskite films. The degraded samples were aged for 9 hours under one sun light illumination (first row) and ion charge deposition (second row), respectively.
Figure 2.11 Topographies (first column) and surface charge density profiles (second column) of untreated (first row), Ion-treated (second row), and light-illuminated (third and fourth row) the perovskite film. The images in the third and fourth row were obtained from light on and off during the measurement, respectively. The images in the third column show 3D plots of topographies colored based on the surface potential values. Both images of light illuminated cases show clear charge trapping
along grain boundaries, but the charge trap is more contrasted when KPFM operation is under the light on
Figure 2.12 High-resolution KPFM images of MAPbI$_3$ film after light soaking and nitrogen cation deposition
Figure 2.13 Time evolution of high-resolution KPFM images and top-view SEM images for the MAPbI$_3$ film. The MAPbI$_3$ films coated on the C$_{60}$(35nm)/ITO substrate were degraded under one sun light illumination in 100% oxygen ambient.
Figure 2.14 Time evolution of KPFM images of MAPbI$_3$ films degraded under different charge generation condition in humidified nitrogen or dry air. (RMS roughness is defined by \[ \text{RMS roughness} = \sqrt{\frac{1}{N} \sum_{i=1}^{N} V_i^2} \] where $V_i$ is the potential difference of pixel i and N is the total number of pixels of the measured area)
Figure 2.15 (a) Normal PL and (b) Time resolved PL measurements of MAPbI$_3$ film before and after positive N$_2$ ion deposition.
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<th>$\tau_3$</th>
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<tr>
<td>MAPbI3_Ref</td>
<td>4.49E-08</td>
<td>1.67E-07</td>
<td>5.71E-09</td>
</tr>
<tr>
<td>MAPbI3_corona</td>
<td>6.89E-09</td>
<td>2.75E-08</td>
<td>7.09E-08</td>
</tr>
</tbody>
</table>

Table 2.1 Time constant of MAPbI$_3$ film from Time resolved PL by 3-exponentials fitting.
Although the easy degradation of hybrid perovskite is a widely known problem, there is no clear theory to date that explains its mechanism from the atomic-scale viewpoint. However, we can readily apply a computational method to simulate how MAPbI$_3$ grain surfaces (including grain boundaries) evolve via our hypothesis on trapped-charge driven degradation. In this study, we first carried out AIMD simulations on the basis of density functional theory (DFT) calculations to examine molecular interaction of gaseous molecules (O$_2$ and H$_2$O) and the evolution of the MAPbI$_3$ grain surface in the presence of charge trapping. In all simulations, we adopted an initial geometry with 2 rigidly fixed MAPbI$_3$ units in the bottom of the crystal, 2 structurally relaxed MAPbI$_3$ units on top of the 2 bottom units to form the surface, and adsorbate molecules on the surface. We bestowed different net charge (+1, 0, or −1) in the unit cell to simulate the trapped charge resulting from the separation of excitons generated by light.

We first carried out simulations for the pure MAPbI$_3$ surface without any adsorbate as a control case, which resulted in no special outcome regardless of the net charge trapped (Figure 2.16). This means that trapped charges themselves do not cause any degradation if there are no gas molecules. To investigate the effect of water molecules in the presence of trapped charge, we performed
simulations for five water molecules randomly placed near the surface (Figure 2.17). In the neutral case with no trapped charge (Fig. 2.17b), all water molecules evolved out of the MAPbI$_3$ surface and clustered with themselves in 2 ps although some of them were quite solidly embedded initially. However, when an electrostatic charge was injected to the crystal, notable changes occurred in atomic motions. In the positively charged crystal (Figure 2.17c), the five water molecules quite evenly placed initially on the MAPbI$_3$ surface started to cluster around a methylammonium cation (CH$_3$NH$_3^+$). The CH$_3$NH$_3^+$ cation ended up effectively solvated by water molecules through hydrogen bonds (N-H···OH$_2$), indicating the dissolution of CH$_3$NH$_3^+$ cations into the surface water on the bulk solid with trapped positive charges. The distance between hydrogen of CH$_3$NH$_3^+$ and oxygen of water during the simulation (Figure 2.18d) obviously shows the existence of hydrogen bonds from which the strong CH$_3$NH$_3^+$-H$_2$O interaction is inferred. Such a dissolution process appears prompted by the weakened bond between CH$_3$NH$_3^+$ cation and the PbI$_6^-$ octahedron unit due to the new electrostatic force field induced by the excess positive charge. In the negatively charged crystal (Figure 2.17), on the other hand, a totally different water-induced degradation pathway was discovered. In contrast to the positively charged crystal, water molecules now aggregate
around an iodide anion ($\Gamma^-$) that is fast moving away from its stable position on the surface due to the electrostatic expulsion by the excess negative charge. The elevated height of $\Gamma^-$ during the negatively charged simulation compared to the neutral or positively charged cases (Figure 2.18e) clearly shows the dissociation of $\Gamma^-$ from the perovskite surface. A subsequent encounter between this hydrated $\Gamma^-$ anion with water molecules would result in volatile species such as hydrogen iodide (HI) and methylamine ($CH_3NH_2$), as shown in the 2 ps snapshot of Figure 2.17, from proton transfer between $CH_3NH_3^+$ and $\Gamma^-$. The calculated bond length of HI at 2 ps in Figure 2.17 is 1.616 Å, which is very similar to the known value of 1.609 Å. It is noted that the generation of these volatile compounds in the presence of water molecules is predicted for the first time by simulation although the release of such volatile compounds from the perovskite solid has been previously suggested [20,44,48]. The deprotonation of organic cation initiated by trapped charge consists of the following three steps as shown in Figure 2.19.

\[
\begin{align*}
I^- + H_2O & \rightarrow HI + OH^- \quad \text{(shown at 1.75 ps of Figure 2.19)} \\
OH^- + H_2O & \rightarrow H_2O + OH^- \quad \text{(shown at 1.80 ps and 1.85 ps of Figure 2.19)} \\
CH_3NH_3^+ + OH^- & \rightarrow CH_3NH_2 + H_2O \quad \text{(shown at 1.90 ps and 1.95 ps)}
\end{align*}
\]
of Figure 2.19)

The overall reaction is seen to initiate from proton transfer from one of the surrounding water molecules to an iodide anion (1.75 ps in Figure 2.19), which is followed by sequential water-to-water proton exchange until it eventually ends up in proton transfer from CH$_3$NH$_3^+$ to OH$^-$ (1.80 ps ~ 1.95 ps in Figure 2.19). Such a sequential proton transfer pathway akin to the Grotthuss mechanism through a series of hydrogen bonds may very well be energetically favorable over direct proton transfer from CH$_3$NH$_3^+$ to I$^-$ due to the stabilization of intermediate and transition states. Our AIMD simulation clearly demonstrates the crucial role of water in the degradation of MAPbI$_3$ crystal in the presence of trapped charges as it enables both the solvation of component ionic species and the sequential proton transfer. To check the effect of unit cell size, we performed additional simulations using 2 × 1 supercell with same charge density (Figure 2.20), and to confirm the effect of time-step and simulation time, we also carried out simulations over 4 ps using 1.5 fs time-step with original unit cell (Figure 2.21). For both kinds of simulations, almost the same events were observed including the solvation of CH$_3$NH$_3^+$ cations with strong hydrogen bonds in positively charged case (Figure 2.20d and Figure 2.21d) and the
encircling of dissociated I\(^{-}\) anion (Figure 2.20e and Figure 2.21e) and CH\(_3\)NH\(_3\)\(^{+}\) cation by water molecules to make a proton transfer reaction in negatively charge case. In addition, the charge sum of 5 H\(_2\)O molecules at the initial geometry (Figure 2.17a) is –0.13, –0.23, and –0.47, and the charge on MAPbI\(_3\) surface is +1.13, 0.23, and –0.53 for positively, neutrally, and negatively charged unit cell respectively, which means that the most of injected charge (+0.90 of +1.00 and –0.76 of –1.00) is delocalized on MAPbI\(_3\) at first.
Figure 2.16 a, Expanded view for the initial geometry of pure MAPbI$_3$ surface. b-d, Temporal snapshots of the AIMD simulated atomic trajectories of MAPbI$_3$ crystal with a charge of b, 0, c, +1, and d, −1. All simulations start with the same initial geometry at 0 ps shown in a. Dotted vertical lines represent the boundaries of actual simulation space, beyond which repeated images of atoms are shown because of the periodic boundary condition.
Figure 2.17 a, Expanded view for the initial geometry of 5 H$_2$O molecules on the MAPbI$_3$ surface. b-d, Temporal snapshots of the AIMD simulated atomic trajectories of MAPbI$_3$ crystal with a charge of b, 0, c, +1, and d, −1. All simulations start with the same initial geometry at 0 ps shown in a. Dotted vertical lines represent the boundaries of actual simulation space, beyond which repeated images of atoms are shown because of the periodic boundary condition.
Figure 2.18 a-c, Snapshots at the 2 ps of the AIMD simulations in Fig. 2.17 with a charge of a, 0, b, +1, and c, −1. d, The distances between hydrogen and oxygen atom shown as navy dotted line in b during the positively charged simulation. e, The relative height of asterisk marked I⁻ in a, b, and c compared to the initial position during each simulation.
Figure 2.19 Temporal snapshots of the AIMD simulated atomic trajectories of 5H$_2$O-covered MAPbI$_3$ crystal with a charge of $-1$ during the proton transfer. Hydrogen bonds in proximity that form proton wires are shown by navy dotted lines. Dotted vertical lines represent the boundaries of actual simulation space, beyond which repeated images of atoms are shown because of the periodic boundary condition.
Figure 2.20 a-c, Snapshots at the 2 ps of the AIMD simulations with a charge of a, 0, b, +2, and c, −2 whose initial geometries are 2 × 1 supercell of Fig. 2.17a. d, The distances between hydrogen and oxygen atom shown as navy dotted line in b during the positively charged simulation. e, The relative height of asterisk marked I− in a, b, and c compared to the initial position during each simulation.
Figure 2.21 a-c, Snapshots at the 4 ps of AIMD simulations with a charge of a, 0, b, +1, and c, −1 using 1.5 fs time-step whose initial geometries are same as Fig. 2a. d, The distances between hydrogen and oxygen atom shown as navy dotted line in b during the positively charged simulation. e, The relative height of asterisk marked $\Gamma^-$ in a, b, and c compared to the initial position during each simulation.
We also investigated the effect of oxygen as it is newly shown to cause complete, rapid degradation of MAPbI$_3$ film following charge deposition (Figure 2.7h) or light soaking (Figure 2.7i). Although Haque and coworkers[24,25] assumed that defect sites of MAPbI$_3$ crystals provide a trap site for oxygen molecules on their way to superoxide formation, we considered no defect sites in our simulation and examined only surface degradation pathways involving oxygen in the presence of trapped charges. As a test case, we first simulated a MAPbI$_3$ surface with four embedded nitrogen molecules with or without trapped charges (Figure 2.22). All embedded nitrogen molecules escaped from the surface and flew away within 2 ps regardless of the net charge, which indicates that nitrogen molecules are not involved in any significant interaction with the constituents of MAPbI$_3$ perovskite, as can be easily expected. We then carried out simulations for oxygen molecules with the same geometry as the nitrogen molecules in our test run. In stark contrast to the nitrogen case, most of oxygen molecules persisted near the surface regardless of the net charge (Figures 2.23b, 2.23c, and 2.23d), which suggests that oxygen is involved in strong chemical interaction with the component species of MAPbI$_3$. Additionally, we observed the formation of oxides in positively (+1) (Figure 2.23c) and negatively (−1) (Figure 2.23d) charged case but
not in the neutral case (Figure 2.23b). Superoxide species were also observed in the presence of a charge, which is in good agreement with previous studies.[25] (See Table 2.2) We had a glimpse of a transient interaction toward oxide formation in the middle of simulation (0.5 ps ~ 1.5 ps in Figure 2.23b), but it did not lead to actual bond formation between oxygen and I\(^-\) anion or Pb\(^+\) cation in the absence of charge. On the other hand, in positively or negatively charged perovskite crystal (Figures. 2.23c and 2.23d), the strong interaction of oxygen with the compositional species of MAPbI\(_3\) resulted in the formation of stable I–O or Pb–O bonds.

To confirm the generation of oxides, we additionally analyzed the I–O, Pb–O, and O–O distances for atoms of interest (those forming the I–O–O–Pb bond at 0.5 ps of Figure 2.23b and 1 ps of Figures. 2.23c and 2.23d) throughout the simulation. In the absence of any charge, the I–O and Pb–O distances are quite short (~2.1 Å and ~2.5 Å, respectively) in the early stage of simulation (0 ~ 1 ps) but become longer past 1 ps (Figures. 2.23e and 2.23f), whereas the O–O distance, to the contrary, start out as a long bond (~1.4 Å) but becomes shorter past 1 ps (Figure 2.23g). These observations mean that oxygen do not form stable bonds with the neighboring Pb or I atoms without charge. On the contrary, in the presence of trapped charge, the I–O, Pb–O, and O–O distances show continued
vibrational oscillation, indicating the existence of permanent chemical bonds (Figures. 2.23e, 2.23f, and 2.23g). To determine the superoxide character of the oxygen atoms in question, we carried out Bader population analysis for the 2-ps geometry with a charge of +1, 0, and −1, and confirmed that only oxygen molecules that are bonded to perovskite (oxygen 1 for +1 charged simulation and oxygen 1 and 2 for −1 charged simulation shown in Figure 2.23) have almost −1 charge (Table 2.2). We additionally examined the vibrational frequencies (or periods) of O–O shown in Figure 2.23g in order to verify the formation of oxides. The vibrational period of the O–O bond with no charge is 24.3(±1.0) fs (at 1 ~ 1.5 ps), which is comparable to the experimental value of 21.4 fs (1,556 cm⁻¹) corresponding to free gaseous oxygen molecule[49], but it becomes much longer (41.1(±8.8) and 39.7(±6.7) fs, respectively) in positively and negatively charged crystal. These longer vibrational periods or lower vibrational frequencies are indicative of a weaker O–O bond and simultaneously stronger I–O and Pb–O bonds in the presence of charge, which should be associated with trapped-charge driven oxidation. Also, the charge sum of 4 O₂ molecules at the initial geometry (Figure 2.23a) was −1.47, −1.65, and −2.07, and the charge on MAPbI₃ surface is +2.47, +1.65, and +1.07 for positively, neutrally, and negatively charged unit cell respectively, which
implies that the most of injected charge (+0.82 of +1.00 and −0.58 of −1.00) is delocalized on MAPbI₃ at first like the preceding water cases. Although quite large portion of injected charge (−0.42) is delocalized on O₂ molecules, it is natural electron transfer process to form a superoxide (O₂⁻) because of the electrophilic property of oxygen.
Figure 2.22  a, Expanded view for the initial geometry of 4N$_2$-covered MAPbI$_3$ surface. b-d, Temporal snapshots of the AIMD simulated atomic trajectories of 4N$_2$-covered MAPbI$_3$ crystal with a charge of b, 0, c, +1, and d, −1. All simulations start with
the same initial geometry at 0 ps shown in a. Dotted vertical lines represent the boundaries of actual simulation space, beyond which repeated images of atoms are shown because of the periodic boundary condition.
Figure 2.23  

a, Expanded view for the initial geometry of 4 O₂ molecules on the MAPbI₃ surface.  

b-d, Temporal snapshots of the AIMD simulated atomic trajectories of MAPbI₃ crystal with a charge of b, 0, c, +1, and d, −1. All simulations start with the same initial geometry at 0 ps shown in a. Dotted vertical lines represent the boundaries of actual simulation space, beyond which repeated images of atoms are shown because of the periodic boundary condition. In the 2-ps snapshots, oxygen molecules are designated by numbers for charge analysis (Table 2.2).  
e-g, Temporal change in the e, I–O, f, Pb–O, and g, O–O distance in the I–O–O–Pb bond.
<table>
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<th>2</th>
<th>3</th>
<th>4</th>
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<td>charge</td>
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<td>nonbonded</td>
<td>−0.08</td>
</tr>
<tr>
<td>−1</td>
<td>status</td>
<td>bonded</td>
<td>−1.04</td>
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<td>−0.98</td>
</tr>
</tbody>
</table>

Table 2.2. Character of chemical interaction with perovskite (top entry) and electrostatic charge (bottom entry in number) of 4 oxygen molecules with the geometry at 2 ps snapshot for 4O$_2^-$-covered MAPbI$_3$ shown in Figure 2.23.
2.3. Conclusions

We found that trapped charges would be responsible for triggering the irreversible degradation in the moisture or oxygen induced degradation of perovskite materials. To verify this, we designed a novel experimental setup enabling the deposition of charges on the surface of perovskites for controlled moisture or oxygen degradation experiments. From this setup, we demonstrated that the perovskite materials degraded irreversibly along grain boundaries only when both moisture or oxygen and trapped charge exist simultaneously. Our study explains both why the degradation begins to occur from the different side of interface between perovskite and charge extraction layer for different charge extraction layers and how light soaking always degrades irreversibly in the presence of moisture or oxygen. KPFM study reveals that charges are trapped preferentially along grain boundaries of perovskites even under uniform deposition of ion charges or uniform illumination of light, which supports our idea of trapped charge driven degradation.

We also investigated interactions between gas molecules and MAPbI₃ crystal in the first 2 ps of perovskite degradation by tracking atomic motions using AIMD simulation. We found that different polarities of trapped charge lead to atomically different routes of degradation. In particular, in the negatively charged crystal, our
calculation for water-covered perovskite clearly showed that the deprotonation of CH$_3$NH$_3^+$ that has so far been only hypothesized in several previous studies can actually occur in water-mediated pathways. On the other hand, in our simulations for oxygen-covered perovskite, we found that trapped charge weakens the Pb–I bond and leads to stronger I–O and Pb–O interactions.
2.4. Experimental methods

Solar cell fabrication
Indium tin oxide-coated (ITO) glass substrates (AMG, 9.5\(\Omega\) cm\(^{-2}\), 25\(\times\)25 mm\(^2\)) were sequentially sonicated in acetone, isopropanol, and deionized water. The cleaned substrate was sufficiently dried in oven in order to eliminate all residual solvents. A 35nm thick \(C_{60}\) layer\([21]\) was densely coated on the ITO glass substrates by using a vacuum thermal evaporator at the constant rate of 0.1Å s\(^{-1}\). For the MAPbI\(_3\) solution, equimolar amounts of methylammonium iodide (MAI), PbI\(_2\), and dimethylsulfoxide (DMSO) (i.e., 159 mg of MAI (Xi’an Polymer Light Technology), 461 mg of PbI\(_2\) (Alfa Aesar), and 78 mg of DMSO (Sigma-Aldrich)) were dissolved in 0.6 ml of dimethylformamide (DMF) (Sigma-Aldrich). After spin coating at 4000 rpm for 20s with ether dripping treatment, the transparent adduct films were sequentially annealed at 65 °C for 1 min and 100 °C for 4 min to form black perovskite films. To prepare the hole trasport material (HTM), 72.3mg of Spiro-MeOTAD(Merk) dissolved in 1ml Chlorobenzen(Sigma-Aldrich) with 28.8 \(\mu\)l of 4-tert-butyl pyridine and 17.5 \(\mu\)l of lithium bis (trifluoromethanesulfonyl) imide (Li-TFSI) solution (520 mg Li-TFSI in 1 ml acetonitrile (Sigma–Aldrich, 99.8%)). The HTM solutions were spin-coated onto the perovskite layer at 2000rpm for 30s. After all process, 50nm gold
(Au) as a counter electrode was deposited on the HTM at the rate of 0.3 Å s⁻¹ by using a vacuum thermal evaporator.

In the case of TiO₂-based devices, a TiO₂ blocking layer was fabricated on FTO-coated glass substrates by spin-coating 0.15 M titanium di-isopropoxide dis(acetylacetonate) (Sigma-Aldrich, 75 wt% in isopropanol) in 1-butanol (Sigma-Aldrich, 99.8%) at sequentially increasing spin rate of 700 rpm for 8 s, 1000 rpm for 10 s and 2000 rpm for 40 s. After spin-coating, the TiO₂ blocking layer was heated at 125 °C for 5 min, and this process was repeated once again. The substrate was annealed at 550 °C for 1 hr. The rest of the processes are identical to the fabrication of C₀₆₀-based devices.

**Characterization**

The cross-sectional and surface images of the perovskite films and the fabricated perovskite solar cells were obtained from a high-resolution scanning electron microscope with a focused ion beam system (Carl Zeiss, AURIGA). The optical absorption spectra of the perovskite films coated on the ITO substrate were measured by UV-vis spectrophotometer (Agilent Technologies, Cary 5000) in the 400-850 nm wavelength range. The XRD patterns of the perovskite films on the ITO glass were collected by using New D8 Advanced (Bruker) in the 2θ range of 5-80 degrees. Photocurrent density-voltage curves were measured by a solar simulator (Oriel Sol3A) with keithley 2400
source meter under AM1.5G, which is calibrated to give 100 mW cm\(^{-2}\) using a standard Si photovoltaic cell (Rc-1000-TC-KG5-N, VLSI Standards). The J-V curves were measured by covering devices with a metal mask having an aperture. (6.76 mm\(^2\)) External quantum efficiency (EQE) was measured by a specially designed EQE system (PV measurement Inc.) with 75 W Xenon lamp (USHIO, Japan) as a source of monochromatic light. The photoluminescence spectra was obtained by a spectrofluorometer(Horiba, Flouoromax-4)

**Experimental setup for ion generation and deposition[35]**

The whole chamber contains two connected chambers: the ion generation(IG) chamber and the ion deposition(ID) chamber. The IG chamber has cylindrical shape with 30 mm diameter and and 35 mm height. It is made of transparent acrylic which makes it possible to see the state of corona discharge during experiment. Stainless steel pin and plate creates highly asymmetric electric field in the chamber, when a bias voltage is applied between the pin and the plate. Gas 1 can flow into the chamber through an inlet on the side wall. The polarity of generated ions is determined by the polarity of applied voltage to the pin. The current of generated ions was measured by Faradaycup electrometer (Keithley Sub-femtoamp romote sourcemeter, 6430). A high-voltage supply (FuG ElektronikGmbH,
HCP140-12500) apply voltage to the pin and the substrate in the ID chamber.

The ions generated in the IG chamber flows through a 115 mm long pipe with a 1.5mm diameter that connects the IG and ID chambers. A tee tube connected to this pipe introduces Gas 2 into the system. The flow rates of Gas 1 and 2 are both controlled by mass flow controllers (MKS instruments, MFC Controller 247D, MFC 1179A). The negatively (postively) charged gas ions are electrostatically attracted and deposited on the positively (negatively) biased substrate with the bias voltage of 2kV (-2kV).

**Corona ion generation**

Nitrogen gas inflow with the flow rate of 2 lpm was transformed into positive nitrogen ions by applying 4.2 kV to the pin. The electric current of generated nitrogen ions was measured to be 20.6-25.8 pA, which is indicative of positive ion generation. To generate negative ions, we used hydrogen gas with the flow rate of 2 lpm and applied negative bias (-1.55 kV) to the pin, which generates the current of -3.2 ~ -4.8 pA. (see Table 2.3)

**Atmosphere control**
To maintain the condition of the air inside the chamber at constant, the flow rate of Gas 2 was controlled by MFC and set as 1.5 lpm throughout the measurements. Nitrogen, dry air, humidified nitrogen and air were used as Gas 2. Nitrogen and hydrogen gases are highly purified by 99.999%, and dry air consists of 80% of nitrogen (99.999%) and 20% of oxygen(99.995%). Gas 2 passes through a water bubbler that controls the humidity in the chamber. The relative humidity was measured by portable multifunction data-logger(Delta OHM, Data logger DO9847, Temp&Humidity probe HP474AC) at the gas exit of deposition chamber.
<table>
<thead>
<tr>
<th>Gas</th>
<th>Gas Type</th>
<th>Flow rate (L m(^{-1}))</th>
<th>Current of generated ion (pA)</th>
<th>Applied voltage</th>
</tr>
</thead>
<tbody>
<tr>
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<td>Nitrogen</td>
<td>2</td>
<td>20.6 ~ 25.8</td>
<td>4.2kV (0.009 mA)</td>
</tr>
<tr>
<td></td>
<td>Hydrogen</td>
<td>2</td>
<td>-3.2 ~ -4.8</td>
<td>-1.55kV (-0.4 mA)</td>
</tr>
<tr>
<td>Gas2</td>
<td>Nitrogen</td>
<td>1.5</td>
<td></td>
<td></td>
</tr>
<tr>
<td></td>
<td>Dry air</td>
<td>1.5</td>
<td>-</td>
<td></td>
</tr>
</tbody>
</table>

Table 2.3 The parameters for ion generation and data obtained from MFC controller and Faraday cup electromer.
Topography and Kelvin probe force microscopy

All the samples for topography and Kelvin probe force microscopy measurements were prepared on ITO glass substrates. The perovskite films were spin-coated on the ITO glass and ITO/C₆₀ substrate. The ion-treated sample was prepared by depositing N₂ positive corona ion for 1 hr on the ITO/perovskite substrate. In the case of the light-illuminated sample, the ITO/C₆₀ substrate was used in order to measure positive charge accumulation profile. After one sun illumination for 1 hr, the sample was measured under light on or off during KPFM operation (Figure 2.11).

Topography and Kelvin probe force microscopy (KPFM) signals were measured by using an atomic force microscope (MFP-3D, Asylum Research) with a Pt-coated tip with the spring constant of 2 nN nm⁻¹ and the resonant frequency of 77 kHz. For each line scanning, topography was first measured and successively the surface potential was measured while scanning the same line at a fixed distance above the sample surface. The tip was positioned sufficiently away from the substrate (as far as 20 nm) in order to exclude the crosstalk artefacts from topography footprints. The surface potential was measured using an active electronic feedback circuitry: the bias voltage to the tip was modulated in order to equate
the potential of the tip with that of the surface, resulting in minimum vibration amplitude of the AFM tip at the fundamental frequency. A 150W halogen bulb with single fiber light guide was used for illumination. (Figures 2.9 and 2.11)

An atomic force microscope (Park systems, NX10) was used to obtain topography and surface potential images of perovskite films. In the ion-treated sample, nitrogen cations produced by corona discharge were deposited on the perovskite film coated on the ITO substrate for 1 hr. In the case of the light-treated sample, the profile of positive charge accumulation in the perovskite film coated on the ITO/C_{60} substrate was measured after one sun illumination for 1 h. (Figures 2.12-2.14)

Computational details

All DFT calculations were performed with the Vienna ab initio simulation package (VASP, version 5.3.5),[50] and the results were visualized by the VESTA (visualization for electronic and structural analysis) and the VMD (Visual Molecular Dynamics) program[51,52]. The projector augmented wave (PAW) method[53,54] was used to describe the electron-ion interaction with the kinetic energy cutoff set to 520.0 eV for the plane waves. We used the Perdew-Burke-Ernzerhof (PBE) exchange-correlation
functional[55] and weak van der Waals interactions were considered by the zero damping DFT-D3 method of Grimme[56]. We performed spin-unpolarized calculations for the neat, water-, and nitrogen-covered MAPbI$_3$ surface but spin-polarized calculations for the oxygen-covered MAPbI$_3$ surface to consider the triplet-state nature of the ground state oxygen. The Brillouin-zone was sampled with a Γ-centered (4 × 1 × 4) Monkhorst-Pack k-point grid for structural relaxation and only a single Γ-point was used for the AIMD simulation. All AIMD simulations were performed in the canonical ensemble using Nosé thermostat with a temperature of 298 K. The total simulation time of each trajectory was 2 ps with a 2 fs time-step. The charge of molecules were obtained by summing atomic charges from Bader population analysis.[57]

The reference structure of MAPbI$_3$ unit cell was constructed from the experimentally determined geometry. We reproduced the experimentally verified orthorhombic MAPbI$_3$ unit cell that consists of four MAPbI$_3$ units (a = 8.836 Å, b = 12.580 Å, c = 8.555 Å)[58], and made little modification on lattice parameters to make the tetragonal crystal structure of MAPbI$_3$ at 25 °C (a = c = 8.852 Å, b = 12.444 Å)[58]. The reference unit cell was fully relaxed at first, whose surface was defined by inserting a large vacuum layer (15 ~ 25 Å) along the b-axis to prevent the interaction with upper periodic
unit cells. To check the effect of vacuum potential, AIMD simulations were performed for neat MAPbI$_3$ surface by changing the vacuum thickness. (Figure 2.23) For the 5 Å case (Figure 2.23a), movable atoms gradually rise up toward the bottom of the upper unit cell because of the potential in the vacuum region. However, for 10 ~ 25 Å cases (Figures 2.23b ~ e), it is clearly shown that no atoms are moved up into the vacuum layer, which implies little interaction between unit cells. Therefore there are negligible dipole effect to simulations. Initial geometries for AIMD simulation were prepared by placing 5 H$_2$O, 4 N$_2$, 4 O$_2$, or a null layer at appropriate positions of the MAPbI$_3$ surface (same locations chosen for N$_2$ and O$_2$). After the initial relaxation in the neutral state, the simulation started with a net charge of +1, 0, or −1 in the unit cell by ejecting (+1) or injecting (−1) an electron. In all surface relaxations and simulations, two MAPbI$_3$ units at the bottom of the unit cell were rigidly fixed while all other atoms that compose the upper two MAPbI$_3$ units and the adsorbates were allowed to move.
Figure 2.24 a, Temporal snapshots of the AIMD simulated atomic trajectories of MAPbI$_3$ crystal in neutral with vacuum thickness of a, 5 Å, b, 10 Å, c, 15 Å, d, 20 Å, and e, 25 Å.
Chapter 3.

Degradation mechanism of other compositional perovskite material

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3.1. Introduction

Organic-inorganic hybrid perovskite materials have lately drawn global attention due to their excellent photovoltaic performance.[2,28,59-62] The organic cations of the material play an important role in photovoltaic properties such as bandgap, exciton binding energy, and carrier lifetime[44,63,64], but they also pose an obstacle in the commercialization of perovskite-based solar cells because of their instability. In general, organic compounds can undergo a chemical reaction with oxygen and atmosphere water, which can be activated under certain conditions such as high temperature or light irradiation. Indeed, the hybrid perovskite materials have shown fast decomposition into organic and inorganic components in the presence of water and oxygen under light irradiation[16,18,19,24,25,30,46,65], which casts serious but
reasonable doubts to the prospect of using such perovskite materials as an active layer in photoelectric devices.

There have been many studies that attempted improving the chemical stability of perovskite materials, mostly by identifying a single source of instability such as water[15,18,19,30], oxygen[16,24,25,46], light[18,65,66], iodine vapor[67] and electric field[18,23]. Recently, we reported a new degradation mechanism based on the electrostatic charges trapped in the perovskite crystal. These charges that were originally generated by light absorption and then trapped mostly along the grain boundaries, were shown to trigger the irreversible degradation of perovskite material in the presence of moisture, which was attributed to the charge driven deprotonation of organic cations. While no similar charge-driven mechanism has been proposed for oxygen-induced degradation of perovskite, we note that the recent discovery of MAPbI$_3$ perovskite suffering from irreversible degradation in the presence of only oxygen without moisture under light soaking[16,24,25,46] indicates that oxygen and water may individually possess different chemical pathways for perovskite degradation.

To dig into this charge driven degradation mechanism, we investigated controlled stability experiments both for commonly used CH$_3$NH$_3$PbI$_3$ (MAPbI$_3$) that is known to form structurally distorted
tetragonal crystals and a new mixed perovskite material having more enhanced structural stability. The crystal structure of perovskite can become more stable by increasing the tolerance factor close to unity by incorporating other organic cation and halide anion with different ion sizes.\[68,69\] Addition of formamidinuim (FA) cation and bromide (Br) anion could not only structurally stabilize the perovskite materials by increasing tolerance factor towards one for inducing cubic crystals, but also enhance their photovoltaic performance by broadening absorption spectrum.\[70,71\] Although several studies on mixed cation and/or halide anion system of \(\text{MA}_x\text{FA}_{1-x}\text{PbI}_y\text{Br}_{3-y}\) have been suggested\[5,70,71\], they focused on the performance aspect, not the stability aspect. Therefore, a new effort is required to develop new composition perovskite ensuring both stability and high performance. In this work, we developed a new mixed \(\text{MA}_{0.6}\text{FA}_{0.4}\text{PbI}_{2.9}\text{Br}_{0.1}\) perovskite ensuring both high performance and stability via lewis-base adduct method.\[3\] Later, it will be shown that this new mixed perovskite still degrades although its degradation speed is slower than the case of conventional MAPbI\(_3\) and the irreversible degradation of both perovskites is triggered by trapped charges.

We also investigated the degradation mechanism of FAPbI\(_3\) which has better absorption property and stability against light and thermal stress than MAPbI\(_3\). A study of the degradation mechanism of FAPbI\(_3\) is lacking
compared to MAPbI₃. Currently, the most anticipated degradation route of FAPbI₃ is the phase transition from black alpha phase to yellow delta phase.\[8,72,73\] Phase transition from alpha phase to delta phase is commonly known to occur under high relative humidity condition at room temperature. In this study, we found that FAPbI₃ also degraded by trapped charges with water or oxygen, but it has different degradation route than MAPbI₃. In the case of pure FAPbI₃ it was confirmed that phase transition to delta phase occurred when there is sufficient water, and we predicted that the reason of phase transition is due to strong interaction between FA organic cation and H₂O molecules. It is also confirmed that the phase is stabilized by adding a little Cs to the FAPbI₃. If the moisture is not sufficient to phase transition, we found that the decomposition of FAPbI₃ occurred by trapped charge.

On the other hand, in the oxygen atmosphere, we found that the decomposition of FAPbI₃ occurred and the phase transition to delta phase hardly occurred. When the decomposition of FAPbI₃ occurs, it finally changes to PbI₂ and the degradation route was predicted by AIMD simulation. In other words, FAPbI₃ is also degraded by trapped charge as MAPbI₃ but predicted to have another degradation mechanism due to the crystal structure and characteristics of FA.
3.2. Results and discussion

In a mixed cation and halide anion system of Pb-based perovskite, APbX₃, the Goldschmidt tolerance factor t is defined in terms of the average radii of ions as follows.

\[ t = \frac{r_A + \bar{r}_X}{\sqrt{2}(r_{Pb} + \bar{r}_X)} \]  \quad (1)

where \( \bar{r}_A \), \( \bar{r}_X \), and \( \bar{r}_{Pb} \) are the average radii of the cation, anion, and lead ions, respectively. The \( I^- \), \( Br^- \) and \( Pb^{2+} \) ions have a spherical shape and the radii of these ions have been measured in previous works.[74] In contrast, it is difficult to precisely estimate the radii of the organic cations perched on the A site due to their non-spherical shape and rotational motion, nevertheless, several groups already reported the calculated effective ion radius of both MA⁺ and FA⁺ cation.[64,75-77] By using equation 1 and the suggested ionic radii \( (r_{MA^+} = 0.18 \text{ nm}, \ r_{I^-} = 0.22 \text{ nm}, \ r_{Pb^{2+}} = 0.12\text{nm}) \), the tolerance factor of MAPbI₃ is calculated to be 0.83[77], indicating that MA⁺ cations are too small to fit into the interstices between PbX₆ octahedra. This mismatch causes crystal distortion, and consequentially MAPbI₃ perovskite have unstable tetragonal crystal structure.[78,79] By partially replacing MA⁺ and \( I^- \) to relatively larger FA⁺ and smaller Br⁻ ions respectively, crystal distortion can be alleviated to produce more stable cubic structure with the tolerance factor between 0.9-
Simple calculation reveals that replacing $\Gamma$ to Br $(r_{Br^-} = 0.196$ nm) only marginally affects the perovskite crystal structure as the tolerance factor of MAPbBr$_3$ is only 0.01 higher than that of MAPbI$_3$. On the other hand, replacing MA to FA cation can significantly alter the crystal structure because FA cations are expected to be much larger than MA cations. The exact radius of FA cation is still controversial,[68] but considering that FAPbI$_3$ can possess not only non-perovskite yellow $\delta$-phase with hexagonal crystal structure ($t > 1$) but also black perovskite $\alpha$-phase with cubic structure ($0.9 < t < 1$), it can be speculated that the tolerance factor of FAPbI$_3$ would be around 1. The radius of FA cation $r_{FA+}$ is estimated to be around 0.26 nm from these speculations and equation 1. The tolerance factor of mixed cation system, MA$_x$FA$_{1-x}$PbI$_3$, is calculated as a function of the ratio $x$ as shown in Figure 3.1. This relation suggests that the $x$ values between 0.2 and 0.6 lead to the most stable cubic crystal structure, as the resulting tolerance factor lies between 0.9 and 1.
Figure 3.1 Dependence of the Goldschmidt tolerance factor on the MA fraction (x) of MA$_x$FA$_{1-x}$PbI$_3$
We measured X-ray diffraction (XRD) patterns of five types of mixed perovskite films (x= 1, 0.8, 0.6, 0.4, 0.2) coated on the ITO glass in order to compare their stability. Figure 3.2a, b show the XRD patterns of the fresh and degraded samples, respectively. The perovskite films were degraded in the chamber (relative humidity (RH) ~ 50%) for 10 hours under one sun illumination. The peak originating from PbI$_2$ at 12.7 degrees intensively appears in the case of MAPbI$_3$ (x=1), and the peak of non-perovskite δ-phase FAPbI$_3$ at 11.7 degrees appears for x < 0.4, which indicates severe instability to water vapor.[8,80] (Figure 3.2c, d) These observations suggest that the most stable composition would be MA$_{0.6}$FA$_{0.4}$PbI$_3$ (x=0.6), which still possesses the tolerance factor ensuring the cubic crystal structure.

We also fabricated the full devices with the ITO/C$_{60}$/Perovskite/Spiro-MeOTAD/Au structure in order to find the optimal composition with respect to the photovoltaic performances in mixed perovskite MA$_x$FA$_{1-x}$PbI$_y$Br$_{3-y}$. Figure 3.3a, b summarize the PCEs as a function of the MA$^+$ fraction(x) and the Br$^-$ fraction(3-y), respectively. As a result, the MA$_{0.6}$FA$_{0.4}$PbI$_{2.9}$Br$_{0.1}$-based devices were shown to the best performance. Considering that MA:FA=0.6:0.4 is the best composition for stability, we concluded that MA$_{0.6}$FA$_{0.4}$PbI$_{2.9}$Br$_{0.1}$ would be the best composition in terms of both performance and stability in mixed perovskite systems.
Figure 3.2 X-ray diffraction patterns of five different mixed perovskite films coated on ITO glass (x=0.2 (black), 0.4 (red), 0.6 (blue), 0.8 (green), and 1 (pink)) (a) before and (b) after degradation under one sun illumination at 50% relative humidity for 10 hours. (c) Magnified XRD patterns around the peaks originating from PbI\textsubscript{2} and non-perovskite δ–phase. (d) Magnified XRD patterns for x=0.2 and 0.4
Figure 3.3 Dependence of the normalized PCEs of ITO/C$_{60}$/Perovskite/Spiro-MeOTAD /Au device on (a) the MA+ fraction (x), and (b) the Br- fraction (3-y)
Figure 3.4a showed the degradation patterns of our mixed MA\textsubscript{0.6}FA\textsubscript{0.4}PbI\textsubscript{2.9}Br\textsubscript{0.1} for two days in the dark condition with relative humidity (RH) 90 %. The absorption spectra of our mixed MA\textsubscript{0.6}FA\textsubscript{0.4}PbI\textsubscript{2.9}Br\textsubscript{0.1} perovskite were hardly changed with the high relative humidity condition (see Figure 3.5) and XRD patterns were the same after two days (see Figure 3.6). This indicates our mixed composition perovskite would be more resistible to become hydrated than distorted tetragonal perovskite MAPbI\textsubscript{3}. It is likely that water molecules could penetrate more easily into the distorted tetragonal MAPbI\textsubscript{3} than into the more compact cubic crystal structure of MA\textsubscript{0.6}FA\textsubscript{0.4}PbI\textsubscript{2.9}Br\textsubscript{0.1}. A slight change of the absorption spectra shown in Figure 3.4b indicates a slow hydration could still happen to our mixed perovskite under 90 % RH. However, MA\textsubscript{0.6}FA\textsubscript{0.4}PbI\textsubscript{2.9}Br\textsubscript{0.1} also perovskites showed irreversible degradation under light soaking at RH 20 % (see Figure 3.4b). It is interesting to note that under light soaking, our mixed perovskite degrades more slowly than the conventional MAPbI\textsubscript{3}.

Moreover, the structurally distorted (conventional) MAPbI\textsubscript{3} film was degraded more quickly than the mixed stable MA\textsubscript{0.6}FA\textsubscript{0.4}PbI\textsubscript{2.9}Br\textsubscript{0.1} under the same moisture and ion deposition level (Figure 3.6).
Figure 3.4 Absorption spectra of (a) $\text{MA}_{0.6}\text{FA}_{0.4}\text{PbI}_{2.9}\text{Br}_{0.1}$ perovskite films under dark conditions at 90% relative humidity. Time evolution of absorption spectra of (b) $\text{MA}_{0.6}\text{FA}_{0.4}\text{PbI}_{2.9}\text{Br}_{0.1}$ during light soaking at 20% relative humidity. Pictures of the perovskite films before and after aging are shown on the left side of each figure.
Figure 3.5 XRD patterns of the perovskite film before (red) and after (black) aged at 90% relative humidity for 2 days.
Figure 3.6 Comparison of the degradation rates of MAPbI$_3$ and MA$_{0.6}$FA$_{0.4}$PbI$_{2.9}$Br$_{0.1}$ perovskite films at 40% relative humidity with positive nitrogen ion deposition. Absorption spectra were measured at an interval of 3 hr.
We examined that the characteristics of MA$_{0.6}$FA$_{0.4}$PbI$_{2.9}$Br$_{0.1}$-based solar cells which has ITO/C$_{60}$/perovskite/Spiro-MeOTAD/Au structure. Figure 3.7a show J-V curves for MA$_{0.6}$FA$_{0.4}$PbI$_{2.9}$Br$_{0.1}$ perovskite solar cells, which demonstrates hysteresis-less performance of the best PCE of 20.2%. The best PCE value was averaged from the J-V curves of forward and reverse scan, which is in agreement with 20.2% of steady-state efficiency shown in Figure 3.7b. The integrated Jsc estimated from external quantum efficiency (EQE) was also well-matched with the measured Jsc as shown in Figure 3.7c. Histograms of the short-circuit current (Jsc), the open-circuit voltage (Voc), the fill factor (FF) and the efficiency of 47 cells are shown in Figures 3.7d, e, f, g. The photovoltaic characteristics of these cells were highly reproducible with a small standard deviation, and the average values are Jsc = 24.34 mA cm$^{-2}$, Voc = 1.058 V, FF = 0.743, and PCE = 19.12%, respectively. This would be the best performance of low-temperature processed perovskite solar cells without using mesoporous TiO$_2$. As seen in Figures 3.4 and 3.6, our new mixed composition perovskite has improved efficiency and stability than MAPbI$_3$, but the MA$_{0.6}$FA$_{0.4}$PbI$_{2.9}$Br$_{0.1}$ also degraded by trapped charge in the presence of water molecules. MA$_{0.6}$FA$_{0.4}$PbI$_{2.9}$Br$_{0.1}$ also decomposed to the PbI$_2$ like MAPbI$_3$.(see Figure 3.8)
Figure 3.7 Photovoltaic performance characteristics. (a) $J$-$V$ curves of the best-performing device measured at various sweep delay times. (b) Stabilized photocurrent density (black) and power conversion efficiency (blue) measured at a bias voltage of 0.91V for 140 seconds. (c) External quantum efficiency (EQE) spectrum and the integrated $J_{sc}$ estimated from the measured EQE. Histograms of (d) short-circuit current density ($J_{sc}$), (e) open-circuit voltage ($V_{oc}$), (f) fill factor (FF), and (g) power conversion efficiency ($PCE$) of 47 cells
Figure 3.8 XRD patterns of the $\text{MA}_{0.6}\text{FA}_{0.4}\text{PbI}_{2.9}\text{Br}_{0.1}$ and $\text{MAPbI}_3$ film before and after degradation by trapped charges in the presence of moisture.
Considering the structural distortion in the MAPbI\textsubscript{3} crystal\cite{69,78,79}, the region between the PbI\textsubscript{6}\textsuperscript{−} octahedron units may be the weakest point of infiltration for external gas molecules. Therefore, structural stabilization of perovskite by tuning tolerance factor to yield a cubic phase structure may benefit its chemical resistance against water or oxygen molecules. In view of the improved stability of our structurally more stable multi-component perovskite MA\textsubscript{0.6}FA\textsubscript{0.4}PbI\textsubscript{2.9}Br\textsubscript{0.1} against water (see Figures 3.4, 3.6 and 3.9), we attempted to check its stability against oxygen as well. We prepared the MA\textsubscript{0.6}FA\textsubscript{0.4}PbI\textsubscript{2.9}Br\textsubscript{0.1} film using our three different experimental apparatuses (Figure 2.7) and measured their absorption spectra and change thereof in 24 hours (Figures 3.10). Surprisingly, these perovskite films showed no sign of degradation under oxygen exposure after 24 hours even with charge deposition and light soaking. To further confirm the improvement in material stability, we measured the photovoltaic conversion efficiency (PCE) of (unencapsulated) devices stored under one sun irradiation in pure oxygen for 36 hours (Figure 3.11 and 3.12), which showed a dramatic improvement over conventional MAPbI\textsubscript{3} based device raising the prospect of these multi-component perovskite materials for long-term stability.
Figure 3.9 Time evolution absorption spectra of MA$_{0.6}$FA$_{0.4}$PbI$_{2.9}$Br$_{0.1}$ (left) and MAPbI$_3$ (right) films aged under one sun illumination in humidified dry air for 18 hours.
Figure 3.10 Absorption spectra of MA$_{0.6}$FA$_{0.4}$PbI$_{2.9}$Br$_{0.1}$ film measured before and after 1 day of aging with dry air, a, under dark condition, b, with nitrogen cation deposition, and c, under one-sun irradiation.
Figure 3.11 Normalized PCEs of MAPbI$_3$ and MA$_{0.6}$FA$_{0.4}$PbI$_{2.9}$Br$_{0.1}$ perovskite solar cell under one-sun irradiation with a continuous flow of pure oxygen. In these devices, C$_{60}$ and Spiro-MeOTAD were used for ETL and HTM, respectively.
Figure 3.12 Time evolution of J-V curves of MA$_{0.6}$FA$_{0.4}$PbI$_{2.9}$Br$_{0.1}$ (left) and MAPbI$_3$ (right) based devices aged under one sun illumination in 100% oxygen ambience for 36 hours.
Next, we investigated the degradation mechanism of FAPbI₃ material. Figure 3.13 shows the degradation patterns of pure FAPbI₃ film in the presence of moisture. First, in Figure 3.13a, 0.5ml of water was placed on the bottom of sealed chamber full filled with humidified nitrogen. The relative humidity of this chamber was over 90%. When pure FAPbI₃ film was placed in this chamber which has lots of water molecules, the black film changed to yellow semi-transparent film. The XRD patterns of yellow film had only δ-phase FAPbI₃ peaks and there were no α-phase FAPbI₃ peaks. (Figure 3.14) That is, in the case of pure α-phase FAPbI₃ under high relative humidity condition, the phase transition occurred faster than the decomposition of FAPbI₃ material occurred. For the experiment under low relative humidity condition (about 35%), we used the apparatus as shown in Figure 2.7. As seen in Figure 3.13b and c, in the case of pure α-phase FAPbI₃ under low relative humidity condition, the decomposition of FAPbI₃ material occurred. The inset figure shows the degraded films at respective condition. When phase transition occurred, the film had very light yellow color. On the other hand, when the FAPbI₃ film was decomposed, the film had dark yellow and a little purple color. This can be seen as PbI₂ and I₂ resulted in decomposition. Thus, when trapped charge is present in FAPbI₃, decomposition occurs in the same as MAPbI₃. However, when moisture is sufficient, the phase
transition occurs differently from MAPbI$_3$. This phase transition can also be predicted by the Goldschmidt tolerance factor. In the case of FAPbI$_3$, the Goldschmidt tolerance factor is more than 1, so it is difficult to maintain the stable cubic structure at room temperature. So AIMD simulation was performed to further analyze the degradation mechanism of FAPbI$_3$. As can be seen in figure 3.15, when the trapped charge was presented, the degradation route of FAPbI$_3$ was almost similar to that of MAPbI$_3$. In positively charged crystal, the FA$^+$ cation was solvated by the interaction between H$_2$O molecules and FA$^+$ cation. As the result of solvation, the FA$^+$ cation was chipped off from the perovskite crystal. In negatively charged crystal, the I$^-$ anion was elevated from the perovskite crystal. Rather, the unusual phenomenon different from MAPbI$_3$ appeared when there was no trapped charges. As mentioned above, in Figure 2.17b, the water molecules evolved out of the MAPbI$_3$ surface in the neutral case. However, the water molecules remained near or inside of FAPbI$_3$ surface especially near the FA$^+$ cation.(Figure 3.15) It means that there is a strong interaction between FA$^+$ cation and H$_2$O molecules. When FA•H$_2$O state is generated by strong interaction, the effective radius of organic cation of perovskite becomes bigger.

Therefore, the Goldschmidt tolerance factor of the FAPbI$_3$ with water molecules is larger than that of FAPbI$_3$ without water, and it is
expected that the phase transition to the $\delta$-phase will occur more easily. The structural instability of FAPbI$_3$ can be solved by doping Cs$^+$ cation.\[[8,9,68]\] Cs-doped FAPbI$_3$ has a small Goldschmidt's tolerance factor and can maintain a stable cubic phase. It was confirmed again in the experimental results. Figure 3.16a is the absorption spectra and picture of the FAPbI$_3$ films aging for one day without light illumination in the ambient air condition (about RH 50%). In the case of pure FAPbI$_3$, the phase transition occurred, but in the case of Cs-doped FAPbI$_3$(5% doping, FA$_{0.95}$Cs$_{0.02}$FAPbI$_3$), there was no changes and the film retained the cubic phase. In addition, the degradation rate of Cs-doped FAPbI$_3$ was much slower than that of pure FAPbI$_3$ under low relative humidity condition.(Figures 3.16b and c) It is suggested that Cs-doped FAPbI$_3$ retains the cubic structure better, thus the penetration of H$_2$O molecules are prevented.
Figure 3.13 Absorption spectra of pure FAPbI$_3$ aged under light illumination at (a) high relative humidity (>90%) condition and (b) low relative humidity condition (about 35%) and positive nitrogen ion deposition at (c) low relative humidity condition. The inset figures are the degraded films under respectively conditions.
Figure 3.14 XRD patterns of FAPbI$_3$ film aged under high relative humidity condition. The inset figures are the picture of film before (α-phase) and after (δ-phase) aging.
Figure 3.15 Temporal snapshots of the AIMD simulated atomic trajectories of 3 H₂O molecules and FAPbI₃ cubic crystal with a charge of +1, 0, and -1. All simulations start with the same initial geometry at 0ps.
Figure 3.16 (a) Absorption spectra of Pure FAPbI$_3$ and Cs-doped FAPbI$_3$ aged in RH 50% air for 1 day. Absorption spectra of Cs-doped FAPbI$_3$ aged under (b) light illumination and (c) positive nitrogen ion deposition at low relative humidity condition.
Figure 3.17 shows the absorption spectra of pure FAPbI₃ aged under 99.995% oxygen. First, we found that the oxygen did not degrade the FAPbI₃ without trapped charges. In Figure 3.17a, the pure FAPbI₃ film was placed in the sealed dark chamber fulfilled with 99.995% oxygen for 2 days. The absorption spectra of FAPbI₃ was hardly changed. Next, we also used the apparatus shown in Figure 2.7b and c for generation charges. As seen in Figures 3.17b and c, the absorption spectra of FAPbI₃ changed rapidly both case of light illumination and positive nitrogen deposition with oxygen. Unlike the case of high relative humidity condition, the phase transition to δ-phase hardly occurred under oxygen condition. The XRD patterns of FAPbI₃ film degraded under light illumination at oxygen condition has only PbI₂ peaks and no δ-phase peaks. (Figure 3.18) It is also confirmed that the degraded film has dark yellow color. From these results, when trapped charges are formed on the perovskite film under oxygen, pure FAPbI₃ was decomposed to PbI₂.

We also carried the AIMD simulations in order to analyze the degradation mechanism of FAPbI₃ with oxygen. (Figure 3.19) In this simulation, the routes of degradation of FAPbI₃ were different from that of MAPbI₃. When there are no charges (Neutral), oxygen molecules escaped from the FAPbI₃ surface. However, when trapped charges existed, oxygen molecules affected to the FAPbI₃ perovskite
crystal. In the case of MAPbI$_3$ with oxygen, it was founded that the strong oxide bonding between Pb-O or I-O was formed. (Figure 2.23) We expected that these generated oxide bond would be loosen the perovskite crystal bonds, as the result the perovskite crystal may have collapsed. On the other hand, in the case of positively charged FAPbI$_3$ crystal, oxygen molecule took place by replacing the FA$^+$ cations. It was expected that the leached FA$^+$ cation was detached from the perovskite crystal, as a result, the FAPbI$_3$ crystals collapsed. In the case of negatively charged FAPbI$_3$ crystal, the strong bond between Pb and O was formed. The bond length between Pb and O was 2.41 Å at 2ps. It was also confirmed that the superoxide was generated in negatively charged case. The oxygen which interact with Pb had -0.62 charges, so it was thought to be super oxide. I$^-$ anion was pushed out from the perovskite crystal because it has -0.73 negative charges. Due to these results, the crystal of FAPbI$_3$ would be destroyed under oxygen condition.
Figure 3.17 Absorption spectra of pure FAPbI$_3$ aged under (a) dark condition, (b) light illumination (c) and positive nitrogen ion deposition at 99.995% oxygen ambient.
Figure 3.18 XRD patterns of FAPbI$_3$ film aged under 99.995% oxygen ambient. The figures are the picture of film before (and after degradation.)
Figure 3.19 Temporal snapshots of the AIMD simulated atomic trajectories of \(2 \text{O}_2\) molecules and FAPbI\(_3\) cubic crystal with a charge of +1, 0, and -1. All simulations start with the same initial geometry at 0ps.
3.3. Conclusion

We have demonstrated that new mixed-perovskite solar cells based on \( \text{FA}_{0.4}\text{MA}_{0.6}\text{PbI}_{2.9}\text{Br}_{0.1} \), whose composition is engineered to minimize the crystal distortion, exhibit improved photovoltaic performance and stability compared to widely-used MAPbI\(_3\) based devices. These cells incorporate C\(_{60}\) as ETL, which is advantageous over TiO\(_2\) due to low-temperature processability, absence of hysteresis, and enhanced device stability. As a result, our best device have shown the PCE of 20.2\%, the highest efficiency of all low-temperature processed perovskite solar cells.

We also demonstrated the degradation mechanism of FAPbI\(_3\) material. In the case of FAPbI\(_3\), the phase transition predominated over decomposition in the presence of water. However, in the case of low relative humidity condition, FAPbI\(_3\) was decomposed by trapped charges. In particular, in AIMD simulation, we found the strong interaction between FA\(^+\) cation and H\(_2\)O molecule which may induces the phase transition. Cs-doped FAPbI\(_3\) was much slowly degraded than pure FAPbI\(_3\) because Cs\(^+\) cation stabilize the cubic structure. In the presence of oxygen, the phase transition of FAPbI\(_3\) was suppressed and the decomposition of FAPbI\(_3\) predominated. In this case, the degradation of FAPbI\(_3\) occurred only when there were
trapped charges. It was predicted that the routes of degradation of FAPbI$_3$ were different from MAPbI$_3$. 
3.4. Experimental method

Solution preparation. All solutions were prepared in the nitrogen-filled glove box. For the MAPbI$_3$ solution, equimolar amounts of methylammonium iodide (MAI), PbI$_2$, and dimethylsulfoxide (DMSO) (i.e., 159 mg of MAI (Xi’an Polymer Light Technology), 461 mg of PbI$_2$ (Alfa Aesar), and 78 mg of DMSO (Sigma-Aldrich)) were dissolved in 0.6 ml of dimethylformamide (DMF) (Sigma-Aldrich). Our multi-component perovskite solution was prepared in the likewise stoichiometric way, i.e., by dissolving 79.5 mg of MAI, 68.8 mg of formamidinium iodide (FAI), 11.2 mg of MABr, 461 mg of PbI$_2$, and equimolar DMSO in 0.6 ml of DMF. FAPbI$_3$ perovskite solution was prepared in the likewise stoichiometric way, i.e., by dissolving 172 mg of FAI, 461 mg of PbI$_2$, and 0.1 ml N-Methyl-2-pyrrolidinone(NMP) in 0.6 ml of DMF. FA$_{0.95}$Cs$_{0.05}$PbI$_3$ was prepared in the likewise stoichiometric way, i.e., by dissolving 168.6 mg of FAI, 13 mg of CsI, 461 mg of PbI$_2$, and 0.1 ml N-Methyl-2-pyrrolidinone(NMP) in 0.6 ml of DMF.

To prepare the hole transport material (HTM) solution, 72.3 mg of Spiro-MeOTAD (Merk), 28.8 μl (26.6 mg) of 4-tert-butyl pyridine (Sigma-Aldrich, 96%) and 17.5 μl of lithium bis(trifluoromethanesulfonyl)imide (Li-TFSI) solution (520 mg Li-
TSFI (Sigma-Aldrich) in 1 ml acetonitrile (Sigma–Aldrich, 99.8%) were dissolved in 1 ml of chlorobenzene (Sigma-Aldrich).

**Perovskite film & device fabrication.** Indium tin oxide (ITO)-coated glass (AMG, 9.5 Ω cm⁻², 25×25 mm²) and pure glass (AMG, 25×25 mm²) substrates were sequentially cleaned by sonication in acetone, isopropanol, and deionized water for 15 min each. All perovskite films were fabricated by Lewis base adduct method[3]. To fabricate a perovskite film, the prepared perovskite solution was coated on a pure glass substrate by spin-coating at 4000 rpm for 20 s with diethyl ether dripping treatment. After spin-coating, the MAPbI₃ film was sequentially annealed at 65 °C for 1 min and 100 °C for 4 min. Our multi-component perovskite film was fabricated via the same method and annealed at 130 °C for 20 min. FAPbI₃ and FA₀.₉₅Cs₀.₀₅PbI₃ perovskite films were fabricated via the same method and annealed at 150 °C for 20 min. To prepare the electron transport layer (ETL), a 35-nm-thick layer of C₆₀ was deposited on the cleaned ITO glass substrate by using a vacuum thermal evaporator (<10⁻⁷ Torr) at the constant rate of 0.1 Å s⁻¹. After the C₆₀ deposition, the perovskite layer was coated by the same process. The HTM solution was then coated on the ITO/C₆₀/perovskite substrate by spin-coating at 2000 rpm for 30 s. Finally, a 50-nm-thick layer of
gold was deposited on the HTM by using the same vacuum thermal evaporator at the constant rate of 0.3 Å s⁻¹.

**Experimental setup for degradation test**

The fully sealed rectangular parallelepiped (RP) chamber has an inner dimension of 35 mm width, 35 mm length, and 30 mm height. This chamber is made of stainless steel to prevent gas leakage except for the top with quartz windows that let light pass through. For ion generation and deposition, two connected chambers, each designated IG for ion generation and ID for ion deposition, were used. The IG chamber produces nitrogen ions by corona discharge that uses an asymmetric field to extract electrons from nitrogen molecules. The ions are then transported by the gas flow and deposited on the perovskite film by the electrostatic force exerted by a high negative voltage applied (−2 kV). The IG chamber is cylindrical in shape (30 mm in diameter, 35 mm in height) and made of transparent acrylic material so that corona discharge can be visually monitored. The ID chamber is also cylindrical in shape (50 mm in diameter, 60 mm in height) and made of stainless steel to block all incoming light. The two chambers are connected to the tee tube so that water molecules or oxygen molecules could be introduced from outside.
The rate of injected gas is controlled by a mass flow controller (MKS Instruments, MFC Controller 247D, MFC 1179A). We used highly purified gases of nitrogen (99.999%) and oxygen (99.995%). Dry air consists of 80% nitrogen and 20% oxygen. To inject water molecules, nitrogen gas was passed through a water bubbler and the relative humidity was monitored by a portable multifunction data-logger (Delta OHM, Data logger DO9847, Temp&Humidity probe HP474AC) at the center of the RP chamber and gas exit of ID chamber. The light soaking experiment was carried out under AM 1.5G one sun irradiation by using a 450 W xenon lamp (Oriel Sol3A) which was calibrated by standard Si photovoltaic cell (Rc-1000-TC-KG5-N, VLSI Standards).

Characterization

Absorption spectra measurements. The absorption spectra of perovskite films were measured by a UV-vis spectrophotometer (Agilent Technologies, Cary 5000) in the wavelength range from 400 to 850 nm.

J-V curve measurement. The photocurrent density-voltage characteristics of perovskite solar cells were measured by a solar
simulator with 450 W xenon lamp (Oriel Sol3A) and a current meter (Keithley 2400) under AM 1.5G one sun irradiation, in which the light source was calibrated by using standard Si photovoltaic cell (Rc-1000-TC-KG5-N, VLSI Standards). A metal mask with an aperture of 7.29 mm² was used in the measurements.

Computational details

All DFT calculations were performed with the Vienna *ab initio* simulation package (VASP, version 5.3.5),[50] and the results were visualized by the VESTA (visualization for electronic and structural analysis) and the VMD (Visual Molecular Dynamics) program[51,52]. The projector augmented wave (PAW) method[53,54] was used to describe the electron-ion interaction with the kinetic energy cutoff set to 520.0 eV for the plane waves. We used the Perdew-Burke-Ernzerhof (PBE) exchange-correlation functional[55] and weak van der Waals interactions were considered by the zero damping DFT-D3 method of Grimme[56]. We performed spin-unpolarized calculations for the neat, water-, and nitrogen-covered FAPbI₃ surface but spin-polarized calculations for the oxygen-covered FAPbI₃ surface to consider the triplet-state nature of the ground state oxygen. The Brillouin-zone was sampled with a Γ-centered (4 × 1 × 4) Monkhorst-Pack k-point grid for structural
relaxation and only a single Γ-point was used for the AIMD simulation. All AIMD simulations were performed in the canonical ensemble using Nosé thermostat with a temperature of 298 K. The total simulation time of each trajectory was 2 ps with a 2 fs time-step. The charge of molecules were obtained by summing atomic charges from Bader population analysis.[57]

We made the cubic crystal structure of FAPbI$_3$ at 25 °C ($a = b = c = 6.36$ Å).[81] Initial geometries for AIMD simulation were prepared by placing 4 H$_2$O, 2 O$_2$, or a null layer at appropriate positions of the FAPbI$_3$ surface. After the initial relaxation in the neutral state, the simulation started with a net charge of +1, 0, or −1 in the unit cell by ejecting (+1) or injecting (−1) an electron. In all surface relaxations and simulations, two FAPbI$_3$ units at the bottom of the unit cell were rigidly fixed while all other atoms that compose the upper two FAPbI$_3$ units and the adsorbates were allowed to move.
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간편 전하에 의한 유무기 페로브스카이트 물질의 비가역적인 분해 현상 연구

서울대학교 대학원 기계항공공학부
곽귀성

요약

유무기 복합 페로브스카이트 결정은 뛰어난 광전특성을 가지고 있어 다양한 광전소자로 활용되고 있다. 특히 페로브스카이트 태양전지는 높은 효율과 저렴한 생산 가격으로 실리콘 태양전지를 대체 할 것으로 예상된다. 이런 페로브스카이트 태양전지의 상용화의 가장 큰 문제점은 낮은 안정성이다. 페로브스카이트 태양전지의 안정성이 부족한 데는 다양한 원인들이 있지만 가장 중요한 원인은 페로브스카이트 물질이 쉽게 분해가 일어난다는 것이다. 페로브스카이트 물질은 공기 중에서 빛을 조사해주면 매우 쉽게 분해가 진행될 뿐 아니라 다양한 요인들에 의해 분해가 진행되게 된다. 따라서 이런 페로브스카이트 물질의 비가역적인 분해현상에 대한 근본적인 원인과 분해과정을 명확하게 밝히는 것이 페로브스카이트 태양전지의 장기 안정성을 확보하는데 반드시 필요하다.

첫째, 페로브스카이트 태양전지에 가장 많이 사용되는 $\text{CH}_3\text{NH}_3\text{PbI}_3$($\text{MAPbI}_3$) 물질의 분해현상에 대해서 연구하였다. 페로브스카이트 태양전지의 경우 전자전달층(Electron transport layer)을 바꿔주는 경우 효율, 하스테레시스, 분해정향성 등이
달라지는 것을 확인하였다. 이는 모두 전하의 기동에 의해서 나타나는 특성이므로 페로브스카이트 분해현상의 원인이 전하일 것이라 추측하였다. 이를 증명하기 위해서 빛이 없는 조건에서 전하를 주입해 주기 위해 코로나 방전법을 이용하였다. 코로나 방전법을 이용하여 주입해준 전하는 Kelvin Probe Force Microscopy와 PL을 이용하여 페로브스카이트 중으로 전달되는 것을 확인할 수 있었다. 이렇게 주입된 전하와 수분 또는 산소가 있는 조건에서 페로브스카이트 물질이 분해되는 것을 확인하였고, 페로브스카이트 중에서 빠져 나오지 못한 간한 전하들이 페로브스카이트 물질의 분해에 핵심적인 원인임을 증명하였다. 보다 이론적인 분해 메커니즘을 분석하기 위해 Density functional theory(DFT) 기반의 ab initio molecular dynamics(AIMD) 시뮬레이션을 진행하였다. 시뮬레이션의 결과 간한 전하가 없으면 아무런 현상도 일어나지 않았으나 간한 전하가 있을때에는 수분 또는 산소와 반응하여 페로브스카이트 결정이 무너지게 된다는 것을 확인할 수 있었다.

둘째, MAPbI₃가 아닌 다른 종류의 페로브스카이트 물질의 탐색 및 분해 현상에 대해 연구하였다. 골드슈미트 팩터(Goldschmidt factor) 계산을 통해 MAPbI₃에서 유기 양이온 또는 할라이드 음이온을 바꿔주면 결정구조가 정방구조에서 입방구조 변할 수 있다. 입방구조의 경우 정방구조보다 구조적 안정성이 뛰어나고 높은 효율을 가질 수 있다. 따라서 CH₃NH₃⁺(MA⁺) 유기 양이온 대신 CH((NH₂)₂)⁺(FA⁺) 유기 양이온을 사용하고 I⁻ 할라이드 음이온 대신 Br⁻ 할라이드 음이온을 사용하여 입방구조의 페로브스카이트 조성을 탐색하였다. 그 결과 20.2%의 효율을 가지면서 MAPbI₃에 비해 뛰어난 안정성을 가지는 MA₀.₆FA₀.₄Pb₁.₉Br₀.₁의 조성을 찾아내었다.
이 조성 역시 간현전하가 있는 경우 분해현상이 나타나는 것을 확인하였으나 특히 산소가 있는 경우 MAPbI\textsubscript{3}에 비해 압도적으로 높은 안정성을 보이는 것을 확인했다.

최근 FA 유기양이온만을 활용하여 입방구조의 FAPbI\textsubscript{3} 물질이 높은 효율과 뛰어난 광 안정성으로 많이 활용되고 있다. 따라서 FAPbI\textsubscript{3}의 분해현상 역시 연구를 진행하였다. FAPbI\textsubscript{3}의 경우 수분이 많은 조건에서는 빠르게 입방구조에서 육방구조로 변하게 되는 것을 확인하였다. 하지만 빛을 조사해 주는 조건에서 산소를 주입하여 주변 결정이 분해되어 PbI\textsubscript{2} 형태로 변하는 것을 확인하였다. 이러한 과정 역시 간현 전하에 의한 효과임을 코로나 방전법과 AIMD 시뮬레이션을 통하여 확인할 수 있었다.

주요어: 페로브스카이트 태양 전지, 유무기 복합 페로브스카이트, 간현 전하, 전자전달층, 분해 메커니즘, 코로나 방전법, AIMD 시뮬레이션, 입방 구조, 혼합 조성 페로브스카이트

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